

Molecular line survey of the high-mass star-forming region NGC 6334I with *Herschel*/HIFI and the Submillimeter Array^{★,★★}

A. Zernickel¹, P. Schilke¹, A. Schmiedeke¹, D. C. Lis³, C. L. Brogan⁴, C. Ceccarelli⁵,
 C. Comito², M. Emprechtinger³, T. R. Hunter⁴, and T. Möller¹

¹ I. Physikalisches Institut der Universität zu Köln, Zùlpicher StraÙe 77, 50937 Köln, Germany
 e-mail: zernickel@ph1.uni-koeln.de

² Max-Planck-Institut für Radioastronomie, Auf dem Hügel 69, 53121 Bonn, Germany

³ California Institute of Technology, MC 301-17, Pasadena, CA 91125 USA

⁴ NRAO, 520 Edgemont Road, Charlottesville, VA 22903, USA

⁵ IPAG, Université Joseph Fourier, CNRS, BP 53, 38041 Grenoble Cedex 9, France

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ABSTRACT

Aims. We aim at deriving the molecular abundances and temperatures of the hot molecular cores in the high-mass star-forming region NGC 6334I and consequently deriving their physical and astrochemical conditions.

Methods. In the framework of the *Herschel* guaranteed time key program CHESS (Chemical *Herschel* Surveys of Star forming regions), NGC 6334I is investigated by using the Heterodyne Instrument for the Far-Infrared (HIFI) aboard the *Herschel* Space Observatory. A spectral line survey is carried out in the frequency range 480–1907 GHz, and further auxiliary interferometric data from the Submillimeter Array (SMA) in the 230 GHz band provide spatial information for disentangling the different physical components contributing to the HIFI spectrum. The spectral lines in the processed *Herschel* data are identified with the aid of former surveys and spectral line catalogs. The observed spectrum is then compared to a simulated synthetic spectrum, assuming local thermal equilibrium, and best fit parameters are derived using a model optimization package.

Results. A total of 46 molecules are identified, with 31 isotopologues, resulting in about 4300 emission and absorption lines. High-energy levels ($E_u > 1000$ K) of the dominant emitter methanol and vibrationally excited HCN ($v_2 = 1$) are detected. The number of unidentified lines remains low with 75, or <2% of the lines detected. The modeling suggests that several spectral features need two or more components to be fitted properly. Other components could be assigned to cold foreground clouds or to outflows, most visible in the SiO and H₂O emission. A chemical variation between the two embedded hot cores is found, with more N-bearing molecules identified in SMA1 and O-bearing molecules in SMA2.

Conclusions. Spectral line surveys give powerful insights into the study of the interstellar medium. Different molecules trace different physical conditions like the inner hot core, the envelope, the outflows or the cold foreground clouds. The derived molecular abundances provide further constraints for astrochemical models.

Key words. submillimeter: ISM – line: identification – astrochemistry – ISM: molecules – ISM: abundances – stars: formation

1. Introduction

The astrochemical composition of the interstellar gas depends on the physical environment and the evolutionary state. In the cold dense regions molecules are unsaturated, while in contrast saturated molecules are much more abundant in hot cores (Herbst & van Dishoeck 2009). Unsaturated molecules are carbon rich with long carbon chains, like radicals as C_nH ($n = 2-8$) or cyanopolynes HC_nN ($n = 3, 5, 7, 9, 11$), whereas saturated organic molecules are hydrogen rich with single bonds of carbon like CH₃OCH₃. Molecular ions are predominantly unsaturated in the cold phase, because most hydrogenation reactions are endothermic or hindered by potential barriers (Herbst 1989). Ices on dust grains begin to form in the cold phase and are observed by broad absorption bands at mid-infrared wavelengths. Sequential hydrogenation of oxygen or nitrogen forms

water or ammonia, respectively. Embedded in the water-ices are relevant amounts of CO, and through further hydrogenation CO → HCO → H₂CO → H₂COH → CH₃OH. Time-dependent astrochemical models suggest that abundances of certain molecules, such as H₂S and SO₂, might be used as an indicator of the evolutionary phase (Nomura & Millar 2004; Herpin et al. 2009), and greater chemical complexity can be expected for longer evolution timescales (Garrod et al. 2008). Approximately 165 molecules have been detected in the ISM or circumstellar shells¹.

Van Dishoeck (2001) emphasized the importance of spectral line surveys for the research of star-forming regions. Line surveys give the possibility to obtain a census of all atoms and molecules and give insights into their thermal excitation conditions and dynamics by studying line intensities and profiles, which allows one to separate different physical components. The discovery of new species or new excited lines from known molecules is possible. Furthermore, the wide range of energy

[★] *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.

^{★★} Appendices are available in electronic form at <http://www.aanda.org>

¹ A current list can be found at <http://www.astro.uni-koeln.de/cdms/molecules>

Table 1. Mass and temperature estimates for NGC 6334I.

Source	$M(M_{\odot})$	T_d (K)	β	n (cm $^{-3}$)	L_{bol} (L_{\odot})
Total ^a	200 ± 100	100	1.7 ± 0.3	1.2×10^7	2.6×10^5
SMA1 ^b	$17 \pm 50\%$	100			
SMA2 ^b	$11 \pm 50\%$	100			
SMA3 ^b	$33 \pm 50\%$	60			
SMA4 ^b	$36 \pm 50\%$	33			

Notes. Values taken from ^(a) Sandell (2000) and ^(b) Hunter et al. (2006). M is the mass, T_d the average dust temperature, β the dust emissivity index, n the hydrogen density and L_{bol} the bolometric luminosity.

levels and eventually opacities of one species allows one a better determination of the temperature and abundance. The uniform recording with the same spectrometer in one observation reduces the calibration uncertainties.

The giant molecular cloud NGC 6334, also known as the “Cat’s Paw Nebula”, lies in the Carina-Sagittarius Arm of the Milky Way at a distance of (1.7 ± 0.3) kpc (Russeil et al. 2012). Studies at infrared wavelengths revealed sites of massive star formation (Persi et al. 2005). NGC 6334I was first explored by Emerson et al. (1973) as the brightest infrared source in the northern part of the cloud, containing an OH maser and an ultracompact H II region. According to Emerson et al. (1973), the pumping of the OH maser requires an energy source of a massive collapsing protostar ($M > 30 M_{\odot}$). H II regions are associated with high-mass stars (spectral classes O and B) where the UV radiation is strong enough to ionize hydrogen. Russeil et al. (2010) studied the densest cores in NGC 6334 in molecular line and dust continuum emission. For NGC 6334I, they derive a total mass of $M = 1206 M_{\odot}$ from the continuum flux density at 1.2 mm, an average density of $n_{\text{H}_2} = 1.2 \times 10^6 \text{ cm}^{-3}$ and an upper limit for the virial mass of $M_{\text{vir}} = 525 M_{\odot}$. The total mass estimate is relatively high, because a low value of the dust temperature of only 20 K was assumed, as compared to the value of $T = 100$ K derived by (Sandell 2000), which leads to a factor of six difference of the mass estimate. While an average dust temperature of $T = 25$ K is appropriate for most cores in the NGC 6334 complex (Matthews et al. 2008), the dust temperature is known to be much higher in NGC 6334I.

Interferometric studies revealed that NGC 6334I consists of several cores, labeled SMA1 to SMA4 (Hunter et al. 2006). The separation distance between SMA1 and SMA2 is about $3.5''$ or 6000 AU. SMA1 and SMA2 show rich spectra of molecular transitions and SMA3 coincides with the H II region which exhibits free-free and dust emission and is excited by the near-infrared source IRS 1E (De Buizer et al. 2002). SMA4 only shows dust emission but no line emission, which is still unexplained. Table 1 summarizes characteristic values for NGC 6334I, derived from the dust continuum emission at different wavelengths, which classify the source as a low-luminosity high-mass hot core. The overall size of NGC 6334I is $10'' \times 8''$, where $1''$ equals to 1700 AU. SMA1 and SMA2 are so-called hot molecular cores, where certain attributes apply according to Cesaroni (2005): a high temperature ($T \geq 100$ K), sizes smaller than 0.1 pc, large masses (10 – $1000 M_{\odot}$) and luminosities exceeding $10^4 L_{\odot}$. Furthermore, they are often associated with water masers and ultracompact H II regions. Hot cores are heated by high-mass protostars which are embedded in a dusty gas envelope. Their rich molecular spectra originate from the evaporation of ice mantles covering the dust grains, which release at $T \sim 100$ K molecules to the gas phase. Molecules released

from grain mantles then begin to drive a fast, high-temperature gas-phase chemistry, forming complex organic species like acetone, beginning from precursors like methanol. Their rotational, vibrational and radiative excitation leads to a characteristic emission spectrum. This explains why the spectrum is poorer in earlier stages like the colder protostellar core NGC 6334I(N) (Brogan et al. 2009; Walsh et al. 2010), $2'$ or 1 pc north of NGC 6334I.

Multiple spectral line surveys have been carried out for NGC 6334I: McCutcheon et al. (2000) 334–348 GHz, Schilke et al. (2006) 459–461 GHz and 817–819 GHz, Thorwirth et al. (2003) 88–115 GHz and 218–267 GHz, Walsh et al. (2010) 84–116 GHz and the latest from Kalinina et al. (2010) in the range 81–242 GHz. Besides, several maser lines have been observed from NH_3 (Beuther et al. 2007), CH_3OH (Walsh et al. 1998) and H_2O (Migenes et al. 1999) as well as bipolar outflows in CO (Laurini et al. 2006; Qiu et al. 2011) and SiO (Gibb et al. 2007). This source is part of the CHESS program (Chemical Herschel Surveys of Star forming regions), a key program of the *Herschel* Space Observatory (Ceccarelli et al. 2010). The aim is to study eight different sites and phases of star formation varying in mass, luminosity, evolutionary state, astrochemical composition etc. by conducting and comparing their spectra. In this article, the entire spectrum (500–1900 GHz) of NGC 6334I is analyzed which was observed by *Herschel* in 2010. Previously, some small sections of the spectrum were studied to find transitions of new molecules not observed before in the interstellar medium (ISM): Oxidaniumyl (H_2O^+) by Ossenkopf et al. (2010) and chloronium (H_2Cl^+) by Lis et al. (2010). A detailed analysis of water has been made by Emprechtinger et al. (2010) and of methylidyne (CH) by van der Wiel et al. (2010).

One difficulty is that if a large beam encompasses a complex, unresolved source structure, the spectrum can become complex due to different overlapping components. The spatial information from the SMA is complementary. The *Herschel* beam is larger than NGC 6334I so that the source is spatially unresolved, while the SMA beam can resolve the individual cores. The SMA maps thus reveal the source morphology and give insights into the distribution of various molecules. This helps to decompose the contributions from different components present in the HIFI spectrum and to differentiate between the two hot cores SMA1 and SMA2, the extended emission and outflows, see Fig. 1.

2. Observations

2.1. Observations with *Herschel*

Observations were conducted with the high spectral resolution receiver system HIFI (Heterodyne Instrument for the Far-Infrared; de Graauw et al. 2010) aboard the *Herschel* Space Observatory (Pilbratt et al. 2010). HIFI is divided into seven bands with different receivers which overlap at the band edges. Each band is separated again into two parts a and b using different polarizations and local oscillator chains. The full frequency range is from 480–1907 GHz, with a gap between 1280 and 1410 GHz. The first five bands employ SIS (superconductor-insulator-superconductor) mixers with an intermediate frequency (IF) bandwidth of 4 GHz. Bands 6 and 7 are based on Hot-Electron Bolometer (HEB) mixers and have an IF bandwidth of 2.4 GHz. The HPBW of the 3.5 m *Herschel* telescope varies with wavelength from $41''$ to $12''$. Table 2 gives a log of the HIFI observations of NGC 6334I and Fig. 2 shows the whole calibrated spectrum. The continuum level increases

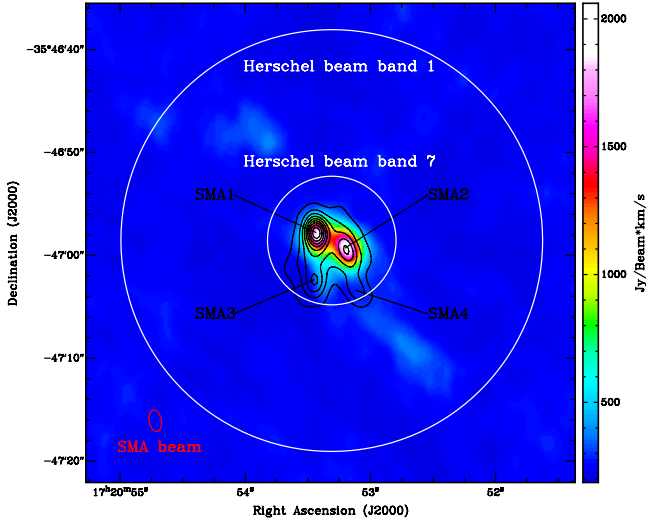


Fig. 1. Integrated SMA line intensity map over all species. Overlaid in black contours is the continuum emission. The contour levels are from 90% to 10% of the maximum in steps of 10%. The large and small white circles represent the FWHM *Herschel* beam size for band 1 and band 7, respectively.

Table 2. Overview over all HIFI bands for the observation of NGC 6334I.

Band	Freq. range (GHz)	Beam size (")	Date	Obs. time (h)	rms (K)
1a	480–560	41	28.02.2010	3.6	0.015
1b	554–636		02.03.2010	2.2	0.03
2a	626–726	29	09.10.2010	3.7	0.04
2b	714–801		19.03.2010	3.3	0.05
3a	799–860	25	09.10.2010	1.8	0.05
3b	858–961		19.03.2010	3.9	0.07
4a	949–1061	21	09.10.2010	4.8	0.09
4b	1052–1122		03.03.2010	2.9	0.10
5a	1108–1244	18.5	13.10.2010	4.2	0.37
5b	1227–1280		12.03.2011	1.7	0.37
6a	1410–1575	14.5	14.10.2010	0.9	
6b	1575–1700		14.10.2010	3.6	1.20
7a	1703–1799	12.5	14.10.2010	2.6	1.27
7b	1788–1907		13.10.2010	3.2	1.27

continuously with frequency, except for band 5b, where a small offset is seen. The line emission in band 1a contributes only 3% to the total flux because the HPBW at 520 GHz is four times bigger than the source size so that in addition the continuum from the surrounding area is observed. The emission lines with the highest peak intensity, $T > 15$ K, are those of CO and ^{13}CO . The telescope was pointed toward the coordinates $\alpha(\text{J2000}): 17^{\text{h}}20^{\text{m}}53.32^{\text{s}}$, $\delta(\text{J2000}): -35^{\circ}46'58.5''$, between SMA1 and SMA2.

The observations were carried out mainly in two sessions in March and October 2010, in dual beam switch mode with an internal chopper mirror switching between the source and two reference positions $3'$ away from the target to provide flat baselines and calibrate the continuum. The total observing time was ~ 42 h. The wide band spectrometer provided a spectral resolution of 1.1 MHz, resulting in a velocity resolution of $0.7\text{--}0.2\text{ km s}^{-1}$. The double sideband spectra were processed with HIPE (*Herschel* Interactive Processing Environment; Ott 2010) version 5, deconvolved to single-sideband spectra following the algorithms described in (Comito & Schilke 2002), and

exported to CLASS (see Sect. 3.1) for subsequent analysis. All spectra are corrected for the beam efficiency and the intensity scale is the main-beam brightness temperature in K. The spectra presented here are equally weighted averages of the horizontal and vertical polarizations, in order to reduce the noise. Band 4a has some sinusoidal baseline features from standing waves. For band 6a only a small portion around a HCl^+ transition was observed, but the data are corrupted and not discussed here. The band coverage is unbiased, meaning that all frequency ranges are included with the band edges where the noise increases. Different sideband gains can lead to discrepancies between line intensities in the two bands in the overlap region, see Fig. B.1g for an example in the $J = 6\text{--}5$ transition of N_2H^+ at 558.966 GHz. The discrepancies amount to 10–20% in the bands 1–3 and become negligible in the higher bands. In the further analysis, this uncertainty can be avoided and has marginal effects when several lines of a given molecule are present.

2.2. Auxiliary observations from the Submillimeter Array

Supporting SMA data at 230 GHz are kindly provided by Brogan and Hunter. The continuum map is identical to that published in Hunter et al. (2006), whereas the spectral image cubes are unpublished. Detailed information about the observation and calibration can be found therein. The data were calibrated again in 2007 in Miriad (described in Brogan et al. 2009). The separation of the line and continuum emission was performed in the uv plane. For the purpose of modeling, the interferometric line and continuum emission were recombined in CLASS for the portion of the spectrum containing the foreground CN absorption feature. The observed frequency range covers two 2 GHz wide bands, 216.6–218.6 GHz in the lower sideband and 226.6–228.6 GHz in the upper sideband, with a frequency resolution of ~ 0.8 MHz. The map covers a $64'' \times 64''$ field and the synthesized beam is $2.2'' \times 1.3''$ in the USB. The line-to-continuum ratio is 12% for SMA1 and 21% for SMA2 in the given bands.

3. Analysis: line identification and modeling

3.1. Simulating spectra with XCLASS

A spectral line is identified by comparing the center frequency with a measured frequency of a known chemical species (atom, molecule, ion) in the laboratory. Two large databases relevant for molecular astrophysics, from the radio to the far-infrared region, have been built up in the last decades: The Cologne Database for Molecular Spectroscopy (CDMS; Müller et al. 2005) and the Jet Propulsion Laboratory (JPL) catalog (Pickett et al. 1998). If the exact source velocity is unknown or several candidates are present for the identification of a spectral line, secondary criteria such as the upper energy level or the Einstein A coefficient must be taken into account. Although many lines in the spectra of hot cores are blended, the identification in spectral surveys is more confident because more lines of a species are present, which can then be checked for consistency, i.e. all lines of a molecule which are expected to be there should be found with the expected line strength, line shape and velocity.

The main program used in this work is XCLASS, which was described in Comito et al. (2005) and applied to a ground-based Orion KL line survey. It is an extension of CLASS (Continuum and Line Analysis Single-dish Software), which is part of the GILDAS² package. XCLASS includes entries from the CDMS,

² <http://www.iram.fr/IRAMFR/GILDAS>

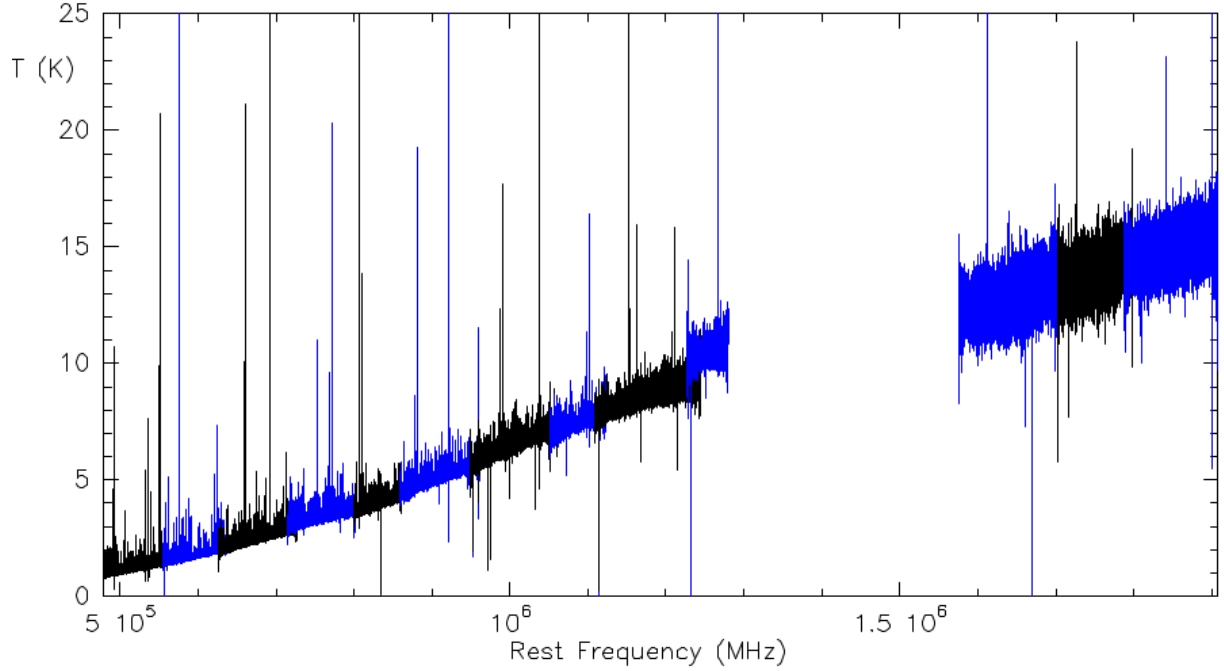


Fig. 2. Complete HIFI spectrum of NGC 6334I. *B*-bands are colored in blue and bands 5–7 are smoothed by averaging over 5–10 channels.

JPL and a private catalog, which contain all necessary information: rest frequency ν_l , integrated intensity, lower state energy E_l , upper state degeneracy g_{up} , quantum numbers and the partition function Q . With this database, it is possible to simulate a spectrum of a molecule, by providing parameters such as the excitation temperature and abundance, and compare the simulated spectrum with the observed one. Similar software tools developed in the last years are *Weeds* (Maret et al. 2011), a package for CLASS, and CASSIS (Caux et al. 2011). The line identification for NGC 6334I, which transfers to a molecule identification in XCLASS, is made easier because 20 detected molecules are reported in the line survey from Kalinina et al. (2010) at lower frequencies and in other surveys mentioned in Sect. 1.

The function which models a spectrum in XCLASS is a solution of the radiative transfer equation for an isothermal object in one dimension, called detection equation (Stahler & Palla 2005). Furthermore, the finite source size and dust attenuation are considered. The model function is:

$$T(\nu) = \sum_m \sum_c \eta(\theta_{m,c}) [J(T_{ex}^{m,c}) - J(T_{bg})] (1 - e^{-\tau(\nu)^{m,c}}) \cdot e^{-\tau_d}, \quad (1)$$

with the optical depth

$$\tau(\nu)^{m,c} = \sum_l \tau(\nu)_l^{m,c} \quad (2)$$

$$\tau(\nu)_l^{m,c} = \frac{c^3}{8\pi\nu^3} A_{ul} N_{tot}^{m,c} \frac{g_l e^{-E_l/kT_{ex}^{m,c}}}{Q(m, T_{ex}^{m,c})} (1 - e^{-h\nu/kT_{ex}^{m,c}}) \phi(\nu)^{m,c}, \quad (3)$$

the dust optical depth

$$\tau_d = N_H \cdot \kappa_{1.3 \text{ mm}} \cdot \left(\frac{\nu}{230 \text{ GHz}} \right)^\beta \cdot 3.3 \times 10^{-26}, \quad (4)$$

the Gaussian line profile

$$\phi(\nu)^{m,c} = \frac{2\sqrt{\ln 2}}{\sqrt{\pi}\Delta\nu^{m,c}} \exp\left(-\frac{(\nu - (\nu_l + \nu_{LSR}^{m,c}))^2}{2\sigma^2}\right) \quad (5)$$

and the beam filling factor

$$\eta(\theta_{m,c}) = \frac{\theta_{m,c}^2}{\theta_{m,c}^2 + \theta_t^2}. \quad (6)$$

The radiation temperature $J(T)$ is defined as $J(T) = \frac{h\nu}{k} \frac{1}{e^{h\nu/kT} - 1}$. The sum goes over the indices m for molecule, c for component and l for spectral line transition. T_{bg} is the background continuum temperature and θ_t the telescope beam size. The line profile function is normalized so that $\int_0^\infty d\nu \phi(\nu) = 1$. The factor $e^{-\tau_d}$ takes the dust attenuation into account where the dust mass opacity κ is taken from Ossenkopf & Henning (1994). For NGC 6334I, we adopt a value of $\beta = 2$ and $\kappa_{1.3 \text{ mm}} = 0.42 \text{ cm}^2 \text{ g}^{-1}$, which is between the values for non-coagulated dust grains with thin ice mantles ($\kappa_{1.3 \text{ mm}} = 0.51 \text{ cm}^2 \text{ g}^{-1}$) and without ones ($\kappa_{1.3 \text{ mm}} = 0.31 \text{ cm}^2 \text{ g}^{-1}$). The total hydrogen column density $N_H = N(\text{H}) + 2N(\text{H}_2)$ is $3 \times 10^{24} \text{ cm}^{-2}$, which is derived from the average hydrogen density given in Table 1 and a source size of $10''$ and is in agreement with values from previous articles (van der Wiel et al. 2010; Lis et al. 2010) considering the uncertainties.

Overall, there are five input parameters for one component of a molecule:

$\theta_{m,c}$	source size in arcsec
T_{ex}	excitation temperature in K
N_{tot}	total column density in cm^{-2}
$\Delta\nu$	velocity width (FWHM) in km s^{-1}
ν_{LSR}	source velocity in km s^{-1} .

Components can be identified as spatially distinct sources within the beam, such as clumps, hot dense cores, colder envelopes or outflows, and can usually be distinguished by different radial velocities. They do not interact with each other radiatively, but are simply superposed in the model. However, for absorption lines the emission from other components is considered first by calculating the background emission spectrum, which is then used as a new “continuum” for absorption lines. By fitting all species

and their components at once, line blending and optical depth effects are taken into account. Because of the high densities in the hot cores ($n_{\text{H}_2} > 10^7 \text{ cm}^{-3}$), it is assumed that most molecules are thermalized at the gas temperature, so that the assumption of LTE is reasonable. With the LTE condition, every population of a level i is known and given by the Boltzmann distribution $n_i = \frac{N}{Q} g_i e^{-E_i/(kT)}$. The modeling of a spectrum speeds up line identification enormously, since several hundred lines of a molecule like methanol are predicted at once in a HIFI band.

The modeling can be done together with isotopologues (and higher vibrational states) of a molecule assuming a conversion factor stored in a separate database. All parameters are expected to be the same except the abundance. For carbon, Milam et al. (2005) derive an average gradient of $^{12}\text{C}/^{13}\text{C} = 6.21(1.00) \cdot D_{\text{GC}} + 18.71(7.37)$, where D_{GC} is the distance to the Galactic center in kpc. The carbon isotopic ratio for NGC 6334I would be 57(10). Eight isotopic species are found in this survey, besides the main isotopes. The ratios used for NGC 6334I are: $[^{35}\text{Cl}]/[^{37}\text{Cl}] = 3$, $[^{32}\text{S}]/[^{34}\text{S}] = 23$, $[^{32}\text{S}]/[^{33}\text{S}] = 127$, $[^{12}\text{C}]/[^{13}\text{C}] = 60$, $[^{14}\text{N}]/[^{15}\text{N}] = 300$, $[^{16}\text{O}]/[^{18}\text{O}] = 500$ and $[^{16}\text{O}]/[^{17}\text{O}] = 1500$. Deuterium is found also, but here no common ratio is taken because the chemical fractionation is very high. The ortho to para ratio for some molecules is assumed to be the statistical equilibrium value (high temperature limit).

3.2. Model fitting with MAGIX

Once an initial guess is set up, the parameters can be fitted using the program MAGIX (Modeling and Analysis Generic Interface for eXternal numerical codes; Bernst et al. 2011; Möller et al. 2012). It delivers a package of different optimization algorithms and an interface to XCLASS. For every of the five fit parameters a range can be set over which to search for the best fit, or the parameter can be kept constant. Used here is the Levenberg-Marquardt (LM) algorithm with the method of least squares, which tries to minimize the squares of the difference between the fit function $y(x_i|a_0 \dots a_{M-1})$ with the M fit parameters a and the data. The sum over all n data points (x_i, y_i) is defined as:

$$\chi^2 = \sum_{i=0}^{n-1} \frac{(y_i - y(x_i|a_0 \dots a_{M-1}))^2}{\sigma_i^2}. \quad (7)$$

The uncertainty for y_i is σ_i . The iteration stops when a minimum is found and the algorithm has converged, or when the change in χ^2 from one iteration step to another is very small, e.g. $\Delta\chi^2 < 10^{-6}$. Since this function is a priori unknown, it becomes clear that the algorithm can converge into a local and not a global minimum. A bad fit for all parameter values can be a hint for non-LTE populations, like HCN lines with high critical densities, or maser activity. The automated fitting has to be corrected manually sometimes, since different problems can occur, which are described in the following.

Starting values: the choice of the initial values for the parameters, is fundamentally important in the LM algorithm. Bad or random starting values lead to an unnecessary amount of computation time. The starting values are set manually by the user and should be physically reasonable. Because of their fit behavior, the five fit parameters in XCLASS are divided into two groups: θ, T_{ex}, N and $\Delta v, v_{\text{LSR}}$. The last two are generally well fixed and do not change anymore, so they are held constant. For high opacities, the line width is broadened which is known as optical depth broadening. To check for a local or global minimum, MAGIX was iterated again with another set of starting values and then

checked if it still converged to the same point. The disadvantage is that this takes computation time with several repetitions and the variation possibilities increase dramatically with more velocity components for a given molecule. This check has been done regularly with molecules with one component, but infrequently for more complicated cases. The best fit parameters for the species with multiple overlapping components may thus not be uniquely determined.

Calculation time: the time increases with the number of parameters and data points. To reduce the calculation time, especially for “weeds”, not all transitions from every band are included but a variety of optical thin and thick lines, lowest and highest frequencies and accordingly a big range of different energy states. Suitable for this are Q-branches of methanol or dimethyl ether. Weeds is the colloquial term for interstellar molecules, which show a plethora of spectral lines. They are all organic asymmetric rotors with a methyl group ($-\text{CH}_3$), except of SO_2 . The classification is not strict, but these five are considered to be the most important ones: methanol, dimethyl ether, ethanol, methyl formate and sulfur dioxide. In methanol, CH_3OH , the methyl group can tunnel the potential barrier at every 120° relatively to the hydroxyl group ($-\text{OH}$), leading to an internal rotation or torsion about the C-O bond. The torsional levels are split into sub-levels (A and E symmetry), and the interaction of torsional and rotational motions yields multiple energy states and a complex spectrum. Furthermore, heavier and larger molecules have a larger moment of inertia, which is inverse to the rotational constants, so that they produce denser spectra. The determination of the excitation temperature with XCLASS is much faster than the so called Boltzmann plot or rotational diagram method, where the upper state abundance N_u of a line transition can be derived from the integrated area in the optical thin limit. This method has the disadvantages that it is unpractical for many lines and misleading for line blends and high optical depths.

Uniqueness: the model function Eq. (1) has two limits where degeneracy occurs: in the optically thick limit $\tau \gg 1$, the last term vanishes and $J(T_{\text{ex}})$ and $\eta(\theta)$ enter linearly into the equation; in the optically thin limit $\tau \ll 1$, $(1 - e^{-\tau}) \approx \tau$ and $\eta(\theta)$ and N enter linearly. This results in the fact that these two parameter pairs cannot be determined independently. In the fitting procedure this can be noticed in the behavior that χ^2 changes only slightly while one of the parameter increases and the other decreases. It is not always obvious if a spectral line is optically thick, especially when several components add together. A solution to the problem of degeneracy are isotopologues, if present. With the ratios mentioned above, a consistent model can be fitted with constraints on N and also for Δv . In some cases the rare isotopologues are fitted well, but the main isotopologue is optically thick and too weak to explain the high observed intensities; this is the case for CH_3OH , SO , H_2CO and HCO^+ . This can be a hint that a second component is present when no signs for that could be deduced from the line shape. However, every additional emission component should be justified, since this makes the model more complex and increases the possibility of redundancy. In the case of NGC 6334I, two near overlapping components are justified by the SMA observations which show that both cores SMA1 and SMA2 show line emission at different velocities.

Weighting: χ^2 is the sum over all modeled spectral lines. One weighting could be done by the uncertainty of y_i , σ_i , which is currently not implemented in MAGIX. But this is not a severe limitation, because the noise or rms is nearly constant across a given HIFI band. It is relevant when the line intensity of the

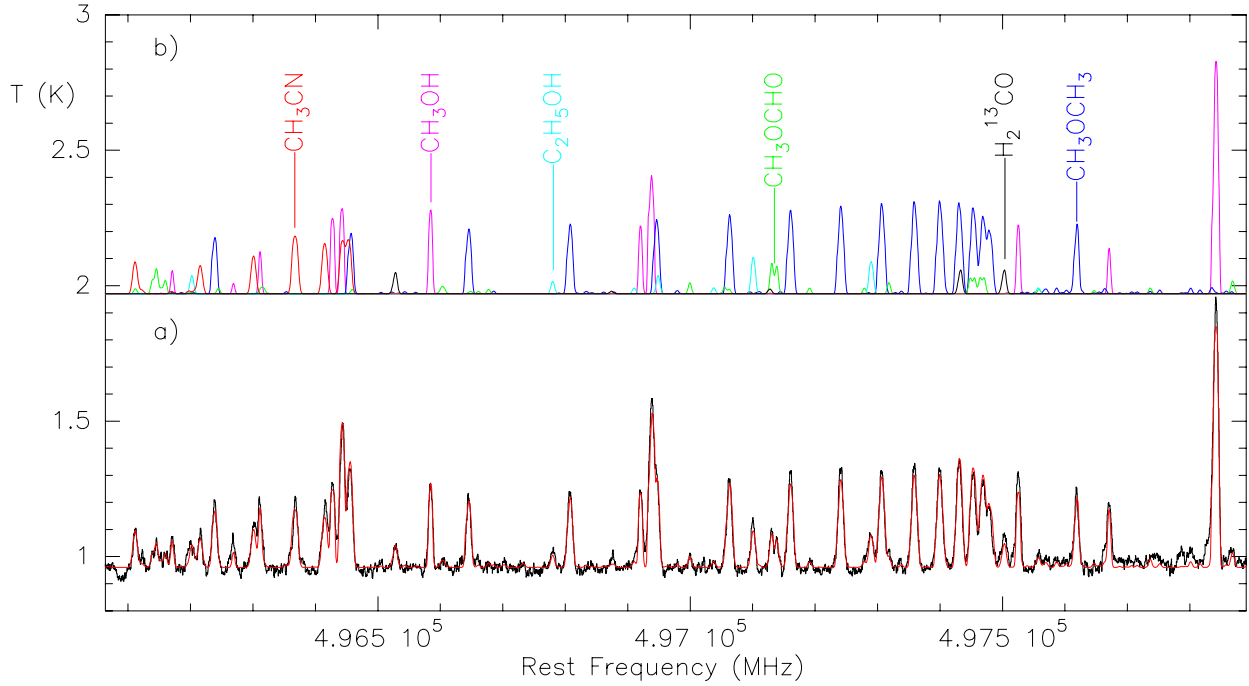


Fig. 3. **a)** Fit example around 497 GHz showing the data (black) and the fit (red). The decomposition into the contributions of individual species is shown above in **b)**.

main isotopologue is much higher than for the less abundant isotopologue. Stronger lines are then more weighted compared to weak lines because the change in χ^2 is larger. This can lead to a model that is biased toward fitting the stronger lines. Another problem is the fitting in higher-frequency bands with larger noise (band 5–7). Here the algorithm tends to increase the line width because it smooths the line wings better. Line widths from transitions in lower bands reveal then that this is an overestimation. To fix this, the line width must be held constant and the best fit value is taken from transitions in lower bands.

Error estimate and confidence limits: it is assumed that in general the standard deviation of θ , T and N is at least 10%, because a small higher or lower value does not change the fit significantly. Additionally, there are other errors not accounted for in the fit like calibration uncertainties, uncertain isotopic ratios, $N(\text{H}_2)$, etc., so that in the worst case only the magnitude of N is roughly known.

Because lines of one species are blended with lines of other molecules, the fitting procedure is done in three major steps:

1. Fitting intense lines and weeds. By modeling weeds one filters them out and takes into account the line blendings. For intense lines of diatomic molecules like SO, the effect of line blending is small. Blended lines are first excluded from the fitting range.
2. Fitting weaker and absorption lines. For weak lines, the noise becomes important so that the fitting procedure is less reliable. Too weak lines, below 2σ , of a molecule are therefore excluded from the fitting range, except for cases where they can serve to give an upper limit for T or N . Absorption lines are not expected to be associated with the hot cores in NGC 6334I because they come mainly from cold foreground gas. Apart from that are ground state transitions with low energy levels from H_2O or NH_3 .
3. Repeating all in a second iteration. The difference now is that the simulated spectrum of all molecules are used as a new background. The parameters are saved globally in a file and read in by MAGIX, but only the molecule of interest is

fitted again. This time, blended lines of a species are taken into account and more components added, if necessary. A simultaneous fit of all parameters of all components of all molecules is in principle possible, but far too time consuming and inefficient. In the end, the result is an improved and consistent fit, see the extract in Fig. 3. This could be done iteratively several times, but it was noticed that the changes are marginal and larger only for molecules with very few transitions present.

3.3. Analysis of interferometric data

The presence of several components with small radial velocity differences which overlap makes the assignment to one certain physical source difficult. The advantage of interferometer data is having the spectral and spatial information together. On the other hand, the SMA frequency range of 4 GHz is small compared to HIFI and extended emission can be missed, because it is filtered out by the interferometer. Two spectra were extracted, one from the peak emission toward SMA1 and the other toward the peak emission of SMA2. It must be noticed that for some molecules like methanol, the emission overlaps from both cores. A direct comparison between them reveals that firstly SMA1 is constantly red-shifted compared to SMA2, secondly the line width in SMA1 is broader and thirdly some lines are only present in SMA1 or SMA2. The intensities in Jy/beam are converted to Kelvin by using the formula

$$T_{\text{mb}}(\text{K}) = \frac{1.22 \times 10^6 \cdot F(\text{Jy})}{[\nu(\text{GHz}) \cdot \theta('')]^2}, \quad (8)$$

where θ is the beam size. The conversion factor is then about 12.5 K/Jy for the LSB and 8.5 K/Jy for the USB. The rms is about 2.5 K in the LSB and 2 K in the USB. Maps of the integrated intensity over one spectral line of all detected molecules were made with the toolkit Karma (Gooch 1996) and the source sizes were derived from Gaussian fits. The whole procedure explained in Sect. 3.2 was performed again with the SMA spectra.

The model derived from the HIFI spectrum was extrapolated to the lower frequencies of the SMA bands and compared with the observed one. Deviations can be explained by the fact that the HIFI model includes the outer parts like the envelope and is an average over a range of temperatures, whereas the SMA spectra correspond to the inner hot core with higher temperatures.

The new derived source sizes θ were adapted for the HIFI spectrum. For some molecules, the source sizes derived from the HIFI spectrum were already in good agreement with the one from the SMA data. Small deviations ($\leq 2''$) are permitted when they lead to a lower χ^2 value. The source size is likely to change when excitation conditions differ, so that the high- J transitions in the HIFI spectra may come from a more compact region than that observed in the SMA maps. For the cases where the lines are optically thin or thick, θ is held fixed to cancel out the ambiguity.

4. Results

4.1. General results

A total of 41 molecules with their isotopologues are identified in the HIFI spectrum. The best fit results are given in Table 3. The whole survey would need several hundred pages of figures and tables with line lists. Therefore, to give an impression, only selected fits for each molecule are attached in Appendix B. The modeled spectrum for each molecule and the full observed and modeled spectrum for NGC 6334I will be made available online, after reprocessing the HIFI spectra with HIPE 8. See Appendix B for further details. The average source velocity (considering all components in the range -5 to -10 km s $^{-1}$ associated with NGC 6334I) is -7.3 km s $^{-1}$ and average line width 4.2 km s $^{-1}$ with a standard deviation of about ± 1 km s $^{-1}$ for both spectral features. The velocity width is dominated by turbulence, since the thermal Doppler width contributes less than 1 km s $^{-1}$ for typical values of $T = 100$ K and $m = 20$ amu. Some molecules show strong broad wings which are produced by outflows: CO, CS, SiO, H₂O and HCN. H₂CS has an asymmetric line profile with a stronger redshifted wing. Especially the component at -17 to -20 km s $^{-1}$, observed in C, NH₃ and H₂S with $T = 10$ – 30 K is presumably the blue outflow lobe. Two cold (~ 5 K), extended foreground clouds are detected consistently with components around -2 km s $^{-1}$ and $+7$ km s $^{-1}$ in the following species: C and C $^+$, CH and CH $^+$, H₂O and H₂O $^+$, HF, H₂Cl $^+$ and OH $^+$. This is in agreement with an observed OH absorption component at $+6$ km s $^{-1}$ toward several NGC 6334 cores reported by Brooks & Whiteoak (2001).

The best fit results for the SMA spectra are presented in Table A.1, the list of detected lines in Table A.2 and the spectra for the LSB and USB in Fig. A.1 in Appendix A. Overall, 20 molecules are identified, of which 15 are also present in the HIFI observations. The additional five molecules are all organic: acetone (CH₃CH₃CO), ethylene glycol ((CH₂OH)₂), formic acid (HCOOH), acrylonitrile (C₂H₃CN) and propionitrile (C₂H₅CN). The latter two only have low intensity features ($< 3\sigma$) in the HIFI band 1a spectrum and are therefore considered to be not detected. Also, 3 isotopologues of OCS are found. In general, the fits give higher temperatures than those derived from the HIFI data, up to 200–300 K, because the emission originates from the center of the hot cores, observed in a smaller beam. Although not true for all molecules, SMA2 has on average higher excitation temperatures, lower column densities, abundances and smaller line widths than SMA1. On average, the source velocity for SMA1 is -5.6 ± 0.1 km s $^{-1}$ and

for SMA2 -7.6 ± 0.1 km s $^{-1}$, so that the relative radial velocity difference is 2 km s $^{-1}$. When applied e.g. to ethanol (C₂H₅OH), which is only seen toward SMA2, the source velocity derived from the SMA spectra is consistent with that from HIFI. In the case of C₂H₄O or NH₂CHO, a small discrepancy of 0.8 km s $^{-1}$ is seen, which can be explained by the fact that in the SMA bands only a few transitions are detected, some of which are blended, so that the precise velocity determination is hindered.

The chemical differences between both cores are determined from the SMA maps and the velocity components in the HIFI spectra. Those molecules that show dominant emission or only emission in SMA2 are HCOOH, (CH₂OH)₂, NH₂CHO, C₂H₄O and C₂H₅OH. For SMA1, the dominant species are HC₃N, C₂H₃CN, C₂H₅CN, H₂S, CN, HCN and HNC. It is noticeable that more oxygen-bearing species, including water (Emprechtinger et al. 2012), arise in SMA2 and more nitrogen-bearing species in SMA1. This variation in the peak abundance is also observed in other star-forming regions, such as Orion KL between the hot core and the compact ridge (Comito et al. 2005).

4.2. Statistics

A spectral line is considered to be detected when the peak intensity exceeds the 3σ level of the random noise. Furthermore, the line width should be comparable to that of all other detected lines. In this way one excludes single channel spikes and sets a lower limit for the integrated line intensity. A lower limit of 2σ would be acceptable when this limit is not taken for a single line of a molecule but rather for many lines since the probability that this is a coincidence decreases when having many lines observed and a good model fit. Nevertheless, some weak features are included in certain cases in the fitting range because they can give useful upper limits for the model.

Overall, 49 chemical species have been found in NGC 6334I so far, listed in Table 4, ranging from fine structure transitions of C, through over mostly di- and triatomic molecules, to organic molecules with 10 atoms. About half of them were found in previous line surveys. Methyl acetylene (CH₃CCH) has been identified by Kalinina et al. (2010), but is too weak in HIFI ($< 1\sigma$). We find 31 isotopologues in our data. The number of lines in all HIFI bands derived from the XCLASS fits is 4217, see Table 5. Hyperfine structure splitting is not considered and counted as one line. For the number of unidentified lines see next section. An automated line counting on the observational spectra gives the same number in orders of magnitude, but is too inaccurate at the band edges and for blended lines. The integrated line intensities are derived from the simulated spectrum, and for bad fits like CO, CS, HCN etc. from Gaussian fits. Because the fits are not perfect, the presented values should be viewed with caution. The line density is 5.5/GHz in the first 5 HIFI bands, ranging from 17/GHz in 1a to < 1 /GHz in 5a. Bands 6 and 7 have a high rms and less than 10 lines each. It is evident that the five weeds mentioned in Sect. 3.2 account for over 75% of all lines, and that methanol is the dominant emitter besides CO. Methyl formate has a very dense spectrum, but is less intense at HIFI frequencies than in both SMA bands where it has the highest number of lines.

4.3. Unidentified features

Table 6 lists the frequencies and strengths of the 75 unidentified lines in the HIFI spectra and the 27 unidentified lines in the SMA spectra. The strength of the lines are derived from

Table 3. Fit results for the HIFI spectrum of NGC 6334I.

Species	Θ ($''$)	T_{ex} (K)	N (cm^{-2})	Δv (km s^{-1})	v_{LSR} (km s^{-1})	Notes	Species	Θ ($''$)	T_{ex} (K)	N (cm^{-2})	Δv (km s^{-1})	v_{LSR} (km s^{-1})	Notes
C	40	56	2.4×10^{18}	4.3	-7.3	b	HCS ⁺	4.0	68	7.6×10^{14}	4.3	-7.1	a
	29	68	2.1×10^{17}	4	0.2		NO	3.5	50	2.7×10^{17}	3.9	-7.5	
	ext	6	2.3×10^{17}	1	-2		SiO	3.4	30	3.0×10^{15}	5	-7.4	
	ext	2.7	4.7×10^{15}	1.3	6.5		HF	ext	2.7	1.2×10^{14}	8.6	-5	
	48	26	4.0×10^{16}	3	-20.6		ext	2.7	1.2×10^{13}	1.9	0.3		
C ⁺	ext	112	1.8×10^{18}	6	-9	1	ext	2.7	2.5×10^{13}	2.2	7		
	ext	4	6.0×10^{17}	3	-0.5		HCl	22	39	1.9×10^{14}	4.3	-6	6
	ext	5	6.0×10^{17}	4.6	6.6		ext	11	1.5×10^{14}	8.6	-10		
CH ⁺	ext	13	1.5×10^{14}	5	-9.8	1	H ₂ Cl ⁺	ext	7	4.9×10^{13}	11	-1.2	7
	ext	8	1.6×10^{14}	6.6	-2		SH ⁺	ext	2.7	2.5×10^{12}	4	-5.7	
	ext	2.7	2.9×10^{13}	4.4	7.8		H ₂ S	3	95	2.5×10^{17}	4	-6.1	8
CH	13.4	95	1.9×10^{14}	2.9	-8.2	1	ext	20	1.0×10^{15}	5	-7.8		
	64	4	6.7×10^{13}	1.9	6.7		18.5	11	8.6×10^{14}	17	-17.5	b	
	38	5	3.0×10^{13}	6	-2.3		CS	2.4	87	3.6×10^{17}	4.5	-6.8	9
NH	39	2.7	6.0×10^{13}	2	-6	1	2.2	73	1.4×10^{17}	3.8	-7.8		
	ext	2.7	9.0×10^{13}	10	-4.6		ext	103	6.5×10^{13}	6	-6.8		
	NH ₂	3.4	108	3.6×10^{15}	4		-8.5	NS	3	89	3.4×10^{15}	5	-7
ext		12	2.9×10^{13}	2	-6.1	OCS	2.5	102	1.2×10^{18}	4	-7.5		
ext		2.7	1.0×10^{14}	30	-13	SO	2.7	70	2.1×10^{17}	4.1	-7.4	10	
NH ₃	8	28	2.3×10^{16}	4.2	-7.7	2	5.8	63	7.6×10^{15}	4.0	-8.0		
	ext	9	6.0×10^{13}	2.9	-6.9		SO ₂	1.8	100	9.5×10^{16}	4.7	-7.8	
	ext	9	2.2×10^{14}	16	-17		H ₂ CS	3.4	106	2.7×10^{16}	5	-7.9	11
NH ₂ D	8	28	2.3×10^{14}	3.8	-8.6	b	CH ₃ OH	2.7	99	1.4×10^{19}	2.9	-7.9	12
	N ₂ H ⁺	31	29	4.3×10^{13}	2.5		-6.8	5.3	61	1.3×10^{17}	2.8	-8.3	
20		35	3.5×10^{13}	2.5	-9.5	CH ₃ OCH ₃	2.1	87	3.0×10^{18}	3.2	-7.4	13	
OH ⁺	ext	2.7	2.0×10^{13}	4	3.3	3	5.0	69	7.0×10^{16}	2.7	-8.1		
	ext	2.7	2.8×10^{13}	6.5	-2		CH ₃ OCHO	3	115	7.1×10^{17}	3.5	-7.65	
H ₂ O	4.2	59	2.1×10^{18}	4.3	-8.3	a	C ₂ H ₅ OH	2	117	1.4×10^{17}	3.9	-8	
HDO	4.4	440	1.2×10^{15}	5.3	-7.7		H ₂ CO	2.7	78	1.5×10^{17}	4.5	-7.4	14
H ₂ O ⁺	ext	24	3.7×10^{13}	8.3	-0.7	4	5.5	52	1.2×10^{16}	2.8	-8.5		
C ₂ H	8.7	32	4.3×10^{15}	5	-6.9		CH ₃ CN	1.9	154	2.8×10^{16}	4.7	-7.3	
CN	14.8	22	1.1×10^{15}	5	-6.4	5	C ₂ H ₄ O	2.1	80	5.4×10^{16}	3	-7.6	
CO	32	60	3.7×10^{19}	4.6	-6.8		H ₂ CCO	3	117	6.0×10^{15}	3	-7.7	a
HCN	3.5	32	2.6×10^{18}	6	-5.6	5	HCO ⁺	5.1	73	1.3×10^{15}	4.7	-7.5	15
	5.0	48	3.0×10^{16}	5	-7.5		31	35	1.4×10^{14}	4.9	-7.2		
DCN	3.5	54	3.7×10^{14}	6	-5.6	a	CH ₂ NH	2.4	75	6.0×10^{15}	5	-7.0	
HNC	3.5	50	1.1×10^{15}	6	-5.6		HNCO	2.0	284	1.5×10^{16}	5	-7.7	
	5.0	40	3.3×10^{14}	4.4	-8.2		NH ₂ CHO	2	117	4.7×10^{15}	4.2	-7.7	
							HC ₃ N	2.8	134	4.2×10^{15}	9	-5.6	

Notes. Several rows for one species denote multiple components. Extended source sizes are abbreviated with ext. (1) Including ¹³CH⁺; (2) including ¹⁵NH₃; (3) only isotopologues H₂¹⁸O, H₂¹⁷O; (4) only isotopologues ¹³CO, C¹⁸O, ¹³C¹⁸O, C¹⁷O; (5) only isotopologues H¹³CN, HC¹⁵N, H¹³C¹⁵N; (6) including H³⁷Cl; (7) including H₂³⁷Cl⁺; (8) including H₂³⁴S; (9) only isotopologues C³⁴S, C³³S, ¹³CS, ¹³C³⁴S; (10) including ³⁴SO, ³³SO; (11) including H₂C³⁴S; (12) including ¹³CH₃OH; (13) including ¹³CH₃OCH₃; (14) including H₂¹³CO; (15) including H¹³CO⁺. (a) τ small. (b) outflow component.

Table 4. List of all detected chemical species in NGC 6334I.

1 atom	2 atoms	3 atoms	4 atoms	5 atoms	6 atoms	7 atoms	8 atoms	9 atoms	10 atoms
C	SO, NO	SO ₂ , HCO ⁺	H ₂ CO	CH ₂ NH	CH ₃ OH	C ₂ H ₄ O	CH ₃ OCHO	CH ₃ OCH ₃	CH ₃ CH ₃ CO
C ⁺	CS, SiO	OCS, H ₂ O	H ₂ CS	HC ₃ N	CH ₃ CN	C ₂ H ₃ CN		C ₂ H ₅ OH	(CH ₂ OH) ₂
	NS, OH ⁺	HCN, NH ₂	HNCO	HCOOH	NH ₂ CHO	CH ₃ CCH		C ₂ H ₅ CN	
	CO, CH ⁺	HNC, HCS ⁺	NH ₃	H ₂ CCO					
	CN, CH	H ₂ S, H ₂ O ⁺							
	HCl, HF	C ₂ H, H ₂ Cl ⁺							
	NH, SH ⁺	N ₂ H ⁺							

Gaussian fits. 211 unidentified features were detected initially, going up from band 1a to 4b, and none in the HEB Bands. All of them are emission lines except of one weak absorption line at 507 686 MHz below 5σ . Most lines are relatively weak. The

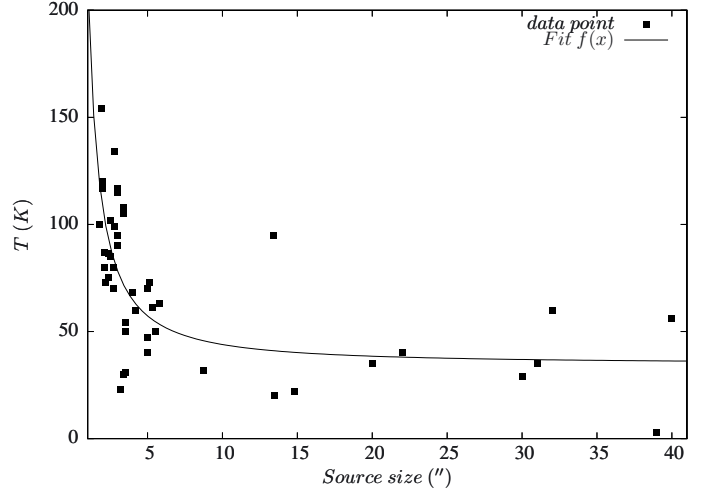
highest peak-to-noise ratio of 14 has the line at 492 325 MHz. Erroneous features are not listed here, but can be checked by comparing the bands in *H* and *V* polarization or at the edges where the bands overlap. In this way, 5 false lines (called ghosts)

Table 5. Integrated spectral line intensities of all detected chemical species of the HIFI spectrum for NGC 6334I.

Species	$\int T dv$ (K km s ⁻¹)	Lines	Species	$\int T dv$ (K km s ⁻¹)	Lines
CH ₃ OH	6016	2063	CH ₃ CN	50	48
CO	4715	31	HCl	40	3
CH ₃ OCH ₃	1486	637	N ₂ H ⁺	23	4
H ₂ O	1257	31	OCS	22	16
CH ₃ OCHO	582	558	H ₂ O ⁺	16	1
H ₂ CO	428	105	NS	16	12
CS	280	45	CN	15	6
C ⁺	262	1	HNC	15	4
C ₂ H ₅ OH	227	134	CH ₂ NH	14	17
NH ₃	226	13	CH	14	1
SO ₂	207	109	OH ⁺	12	2
HCN	205	27	NO	11	8
NH ₂	181	6	H ₂ Cl ⁺	11	2
H ₂ S	168	32	C ₂ H	11	6
HCO ⁺	152	10	HNCO	10	25
HF	146	1	H ₂ CCO	4	9
H ₂ CS	141	78	HCS ⁺	4	4
SO	119	51	NH ₂ CHO	2	8
C	115	2	HC ₃ N	2	4
CH ⁺	108	3	SiO	2	1
NH	103	3	SH ⁺	1	1
C ₂ H ₄ O	67	94			
Sum	17 485	4217			

Notes. Included are all transitions from isotopologues (except CH₃OD and CH₃¹⁸OH) and torsionally or vibrationally excited states.

at 720 GHz are noticed which are present in 2a, but not 2b. They arose presumably from a feedback in the LO and are no longer present after calibration with a newer version of HIPE. Notable are patterns like branches where several lines (5–10) appear at once in a narrow frequency range (1 GHz), namely at 510, 532, 537, 602, 641, 695, 720, 1009 and 1069 GHz. Other clear patterns like a recurring difference between lines could not be discovered. If there are several candidates for an unidentified line, the simulation with XCLASS helps to constrain them. Any anticoincidence, which means that a simulated line in other quantum numbers in other frequency bands does not match with the observed spectrum, is an indication for a false identification. That is why complete line surveys are better suited for confident line identifications than small windows or single line searches. Another constraint from interferometric data is that all lines of a species originate from the same source. Since weeds have the highest number of lines, it is most probable that the unidentified lines originate from them, or one of their isotopologues not yet included in the database. Two papers were checked with reported line frequencies of deuterated methanol CH₃OD (Walsh et al. 2000) and the isotopologue CH₃¹⁸OH (Fisher et al. 2007). 8 lines could be assigned to deuterated methanol and 128, especially most of the former reported branches, to CH₃¹⁸OH. Subtracting them, the remaining number of unidentified lines is 75 lines (including the 1009 GHz and 1069 GHz branches) or about 2% of the total of 4428 lines. In the SMA bands, 144 transitions are counted of which 27 are unidentified. That makes 19% and is much more than in HIFI frequency range. It is probable that they originate from organic molecules in higher excited states. All unidentified lines in the SMA spectra are detected in both cores, but it is remarkable that, with the exception of two, they are all stronger in SMA2 than SMA1.

**Fig. 4.** Derived excitation temperatures in dependence of the source size for molecules associated with NGC 6334I.

5. Discussion

5.1. General remarks

Schilke et al. (2006) analyzed spectra of NGC 6334I with APEX at 460 and 810 GHz with a bandwidth of 2.5 GHz. 12 molecules were identified and roughly fitted, of which 6 show an acceptable agreement (same order of magnitude in N) with our results: CH₃CN, CH₃OCH₃, CH₃OCHO, HNCO, N₂H⁺ and SO₂. The deviation for other molecules can be explained by an over- or underestimation of the source size and by setting the temperature in steps of 50 K. In general, there is a good agreement between the HIFI and SMA column densities. The temperatures given in Table A.1 are not well determined, because for many molecules only one transition is present.

Because no error bars are given, it should be pointed out that the given parameters for some molecules are less reliable than for others. The spectra of strong lines (CO, N₂H⁺, HCN, C⁺...) could be affected by emission in the off-beam position. In some cases, the assumption of LTE does not hold, e.g. for HCN, or the ambiguity leads to an uncertainty. Reliable values are those which are checked consistently with the SMA and have isotopologues, so that better constraints are available for θ and N .

In Fig. 4, the excitation temperature is plotted as a function of the source size for all molecules associated with NGC 6334I. The extended source sizes ($\theta \gg$ beam size) are not included. The data points can be roughly fitted by a function $f(x) = a \cdot x^b + c$ with an exponent of $b = -1.3 \pm 0.5$ and a constant $c = (35 \pm 11)$ K for $\theta > 6''$ in the outer parts or envelope where the temperature gradient becomes flat. This is consistent with models of hot cores which predict a steep increase inwards (Rolfes et al. 2011). The profile could be affected by the fit function, where in the optical thick limit, T and θ are inversely proportional to each other with $T \propto 1/\theta$, but this only concerns few molecules.

The fractional abundances $X = N_{\text{Species}}/N_{\text{H}}$ are distributed widely, ranging over eight orders of magnitude from the nitrogen hydrides ($X = 10^{-12}$) to CO ($X = 10^{-5}$, derived from the isotopologues). The hydrogen column density for NGC 6334I is $N_{\text{H}} = 3 \times 10^{24}$ cm⁻², see Sect. 3.1. The most abundant molecules with the highest column densities after CO are CH₃OH and CH₃OCH₃ (both $X = 10^{-6}$), HCN and H₂O (both $X = 10^{-7}$). The hydrogen column density can be estimated from the derived CO column density of $N = 3.7 \times 10^{19}$ cm⁻² by using a standard fractional abundance of $[\text{CO}]/[\text{H}] = 5 \times 10^{-5}$.

Table 6. List of unidentified lines in the SMA and HIFI spectra for NGC 6334I.

Frequency (MHz)	Area (K km s ⁻¹)	<i>I</i> (K)	<i>I</i> /rms	Width (km s ⁻¹)	Frequency (MHz)	Area (K km s ⁻¹)	<i>I</i> (K)	<i>I</i> /rms	Width (km s ⁻¹)
SMA									
216 671	41.2	8.1	3.2	4.8	227 128	68.0	11.3	5.7	5.7
216 715	71.7 ^b	13.5	5.4	5.0	227 239	26.0 ^b	8.1	4.1	3.0
216 913	48.0	12.4	5.0	3.6	227 244	32.3 ^b	7.0	3.5	4.3
216 919	82.8	24.1	9.6	3.2	227 534	61.0	14.2	7.1	4.0
217 000	54.9	11.4	4.6	4.5	227 728	43.7	12.8	6.4	3.2
217 056	90.0	20.8	8.3	4.1	227 803	41.6	11.4	5.7	3.4
217 262	91.5 ^b	22.3	8.9	3.9	227 823	47.8 ^b	7.4	3.7	6.1
217 427	38.2	11.6	4.6	3.1	227 841	99.0 ^b	21.1	10.6	4.4
217 496	70.9	18.8	7.5	3.5	228 001	53.0	7.6	3.8	6.6
217 593	88.3	24.5	9.8	3.3	228 141	112 ^b	17.0	8.5	6.2
217 678	43.5	9.6	3.8	4.3	228 230	24.5	7.5	3.8	3.1
217 689	119	27.4	11.0	4.1					
217 710	120	27.7	11.1	4.1					
217 738	76.1	18.0	7.2	4.0					
218 158	116 ^b	18.8	7.5	5.8					
218 600	58.8	11.8	4.7	4.7					
HIFI									
492 313	0.88	0.30	14.3	2.8	1004 810	2.13	0.50	4.2	4.0
501 501	0.40 ^b	0.06	4.0	6.3	1007 185	2.50	0.44	3.1	5.3
504 512	0.14 ^b	0.05	3.3	3.0	1008 602	2.10	0.33	2.5	6.1
505 339	0.62	0.10	6.3	6.2	1008 962	3.16	0.47	3.4	6.3
507 167	0.72 ^b	0.07	4.7	9.8	1009 245	1.61	0.34	2.4	4.4
507 673	0.33	0.07	4.3	4.8	1009 460	1.85	0.54	4.2	3.2
515 619	0.25	0.06	4.3	3.6	1009 618	2.31	0.62	5.6	3.5
519 093	0.70 ^b	0.043	5	9.1	1009 728	2.78	0.51	4.3	5.1
528 091	0.33	0.06	4.1	5.0	1009 804	2.90	0.53	4.4	5.1
530 769	0.28 ^b	0.09	6.0	2.8	1009 850	2.74	0.58	4.8	4.5
531 303	0.36 ^b	0.062	4.1	5.4	1009 881	3.20	0.50	4.2	6.0
531 386	0.65	0.09	6.0	6.8	1017 993	4.00	0.71	7.6	5.2
534 943	0.63 ^b	0.049	3.3	12	1024 147	2.78	0.60	6.0	4.3
536 632	0.14	0.06	3.7	2.4	1040 932	1.86	0.52	3.7	3.3
539 134	0.36	0.08	5.3	4.2	1045 510	2.30	0.54	6.0	4.0
556 311	0.44	0.10	6.7	4.1	1047 493	2.00	0.47	3.6	4.1
602 910	0.53	0.10	3.1	4.8	1047 982	3.34	0.62	4.1	5.0
649 022	0.51	0.12	4.0	4.0	1052 768	4.50 ^b	0.75	3.0	5.7
649 093	0.47	0.12	4.0	3.7	1065 139	3.90	0.60	4.6	6.2
650 170	0.74	0.13	3.8	5.4	1066 454	2.00	0.49	4.9	3.8
692 016	0.39	0.10	2.5	3.7	1069 053	1.45	0.30	3.0	4.5
693 184	0.60	0.13	3.3	4.4	1069 207	1.00	0.32	3.2	2.9
695 857	0.41	0.13	3.3	2.9	1069 495	1.35	0.34	3.4	3.7
697 351	0.95	0.12	3.0	7.7	1069 625	1.35	0.32	3.2	3.9
698 919	1.60 ^b	0.27	6.8	5.6	1069 741	1.48	0.34	3.4	4.1
706 552	0.55	0.16	4.0	3.2	1069 851	2.70 ^b	0.43	3.9	5.9
717 678	0.74	0.24	5.3	2.9	1076 527	3.80	0.85	7.7	4.2
741 290	1.24	0.15	3.0	7.9	1078 549	2.63	0.51	4.6	4.8
743 318	0.81	0.15	3.0	5.2	1088 059	2.00	0.55	4.6	3.4
781 479	1.40	0.22	4.4	6.0	1088 302	1.65	0.44	3.7	3.4
785 042	0.85	0.18	3.6	4.4	1095 012	2.86	0.69	6.3	3.9
790 927	3.30	0.52	5.8	6.0	1095 149	1.70	0.45	4.1	3.5
791 560	1.00	0.27	3.4	3.4	1098 546	2.15	0.45	4.5	4.4
795 572	1.60	0.26	3.7	5.7	1099 864	2.90	0.62	6.2	4.4
840 812	0.80	0.21	4.2	3.6	1099 903	1.50	0.32	3.2	4.4
885 659	1.05	0.31	4.4	3.2	1108 349	1.96	0.51	3.6	3.6
972 428	1.70	0.40	3.6	4.0	1112 291	4.00	0.48	3.7	7.7
1 000 929	3.90	0.67	7.4	5.4					

Notes. Given are the center frequency in the velocity rest frame of SMA2, the line area, the peak intensity, the corresponding signal-to-noise ratio, and the line width (FWHM). ^(b) blended line.

The result is $N_{\text{H}} = 7.4 \times 10^{23}$ for a source size of $\theta = 32''$ and is lower compared to the value given above where $\theta = 10''$ is assumed. The following differences and similarities are evident when comparing NGC 6334I with an example of a

low-mass protostar IRAS 16293-2422 and the high-mass star-forming region Orion KL, Table 7: The methanol abundance $X = 5 \times 10^{-6}$ in NGC 6334I is higher than in the Orion HC and IRAS 16293, which have the same ratio. As is shown

Table 7. Comparison of the fractional abundances in NGC 6334I with IRAS 16293-2422 and the hot core (HC) and compact ridge (CR) of Orion KL.

Molecule	NGC 6334I	IRAS 16293-2422 ^c	Orion KL HC ^e	Orion KL CR ^e
CO	1.2(−5)	4.0(−5) ^d	9.0(−5)	1.1(−4)
CH ₃ OH	4.7(−6)	3.0(−7)	1.4(−7)	4.0(−7)
CH ₃ OCH ₃	1.0(−6)	2.4(−7)	8.0(−9)	1.9(−8)
HCN	8.7(−7)	–	3.0(−7)	–
H ₂ O	6.7(−7)	–	5.3(−6)	1.0(−4)
OCS	4.0(−7)	2.5(−7)	1.1(−8)	3.0(−8)
CH ₃ OCHO	2.4(−7)	2.0(−7)	1.4(−8)	3.0(−8)
CS	1.2(−7)	3.0(−9) ^d	6.0(−9)	1.0(−8)
SO	8.3(−8)	2.5(−7)	1.9(−7)	3.0(−7)
H ₂ S	8.3(−8)	9.0(−8)	2.7(−6) ^f	–
H ₂ CO	5.0(−8)	6.0(−8)	7.0(−9)	4.0(−8)
SO ₂	3.3(−8)	1.0(−7)	1.2(−7)	1.6(−7)
HCOOH	1.8(−8)	6.2(−8)	8.0(−10)	1.4(−9)
C ₂ H ₅ CN	1.4(−8)	1.2(−8)	3.0(−9)	5.0(−9)
H ₂ CS	1.0(−8)	5.5(−9)	8.0(−10)	1.2(−9)
CH ₃ CN	9.3(−9)	1.0(−8)	4.0(−9)	5.0(−9)
HNCO	5.0(−9)	9.0(−9)	5.8(−9)	–
C ₂ H	1.4(−9)	2.2(−10) ^d	5.5(−10)	5.3(−9)
HC ₃ N	1.4(−9)	1.0(−9)	1.8(−9)	6.0(−9)
SiO	1.0(−9)	4.5(−9)	6.0(−9)	8.0(−9)
HCO ⁺	4.3(−10)	1.4(−9) ^d	1.1(−9)	1.0(−9)
CN	3.7(−10)	8.0(−11) ^d	8.0(−10)	8.0(−10)

Notes. Notation a(b) indicates $a \times 10^{-b}$. ^(d) Referring to the cooler, less dense envelope.

References. ^(c) Cazaux et al. (2003), Schöier et al. (2002). ^(e) Sutton et al. (1995), Caselli et al. (1993) and references therein. ^(f) From Persson et al. (2007).

by time-dependent chemical models of hot cores (Nomura & Millar 2004), the parent molecule methanol (formed on grain surfaces) is stable for timescales of at least 10^4 yr. The other organic molecules, dimethyl ether and methyl formate, are more abundant in NGC 6334I than in Orion, reflecting perhaps that NGC 6334I is more time evolved. These second generation molecules start to form in the gas phase between evaporated molecules and peak up in abundance at 10^{4-5} yr. After that, the parent molecules are eventually destroyed. For sulfur-bearing molecules, an anticorrelation between NGC 6334I and Orion is observable: while H₂S is two orders of magnitude less abundant in NGC 6334I, OCS is much more abundant than in Orion. The higher depletion of H₂S could be explained by the fact that this molecule is efficiently destroyed by H₃O⁺ in hot cores and subsequently transformed into SO and SO₂ (Wakelam et al. 2004). Overall, the fractional abundance of the sulfur-bearing molecules and their sequence in NGC 6334I is more comparable to IRAS 16293 than to Orion (except for CS). Some nitrogen-bearing molecules like nitriles have about the same fractional abundances in NGC 6334I as in Orion and IRAS 16293: CH₃CN with $X = 1 \times 10^{-8}$, CN with $X = 4 \times 10^{-10}$, HC₃N with $X = 2 \times 10^{-9}$ and C₂H₅CN with $X = 1 \times 10^{-8}$. When comparing fractional abundances in the literature, it should be pointed out that possible uncertainties are the adopted H₂ column density and the source size. It needs to be further investigated if all these deviations are due to different precursor abundances, varying gas temperatures and densities or are due to age effects solely.

Results for most molecules are discussed briefly below, see Figs. B.1 and B.2 for their spectra. They are grouped by their moment of inertia, ranging from the lightest diatomic to the heaviest complex organic molecules. The electronic state for

molecules is denoted by $^{2S+1}\Lambda_{\Omega}$, where S is the total electronic spin, Λ the projection of the total electronic orbital angular momenta along the internuclear axis and $\Omega = \Lambda + \Sigma$, where Σ is the projection of S onto the axis.

5.2. Atomic and diatomic species

C I and C II, Fig. B.1a. Two fine structure transitions of neutral carbon fall into the submillimeter range, $^3P_1-^3P_0$ at 492 161 MHz and $^3P_2-^3P_1$ at 809 342 MHz. Up to five velocity components can be observed, with the main emission at -7.3 km s^{−1}. The nondetection of ¹³C is used to set an upper limit for the abundance. The [C II] $^2P_{3/2}-^2P_{1/2}$ fine-structure transition at 1.9 THz is one of the most important cooling lines in the ISM and one of the most intense single lines ($\int T dv = 262$ K km s^{−1}) in the HIFI spectrum besides CO. Three components are found, corresponding to the ones in C I. The large source sizes for all components indicate a broad distribution. The dust attenuation is neglected in the fit of the C II emission line because it is assumed that most of the emission originates from the diffuse outer layer. Therefore, the derived values for N and T should be taken as lower limits, and we notice that some of the spectral features could be affected or caused by contamination in the off-beam.

HF, Fig. B.1s. The ground transition $J = 1-0$ of hydrogen fluoride at 1 232 476 MHz is seen in absorption (Emprechtinger et al. 2012). Three components at -5 , 0.3 and 7 km s^{−1} are visible. The broad velocity width of the -5 km s^{−1} component indicates that HF is widely distributed and extended over NGC 6334I. HF is considered to be a good tracer of H₂ (Neufeld et al. 2005), since the formation is directly coupled through the exothermic reaction $F + H_2 \rightarrow HF + H$, and the ratio in the diffuse medium is $X = 3.6 \times 10^{-8}$. This would correspond to $N_H = 3 \times 10^{21}$ cm^{−2} in this diffuse layer.

SO, Fig. B.2d. Sulfur monoxide is a radical with two unpaired electrons (ground state $^3\Sigma$) and a spin-multiplicity of 3, so that each total angular momentum level $J = N + S$ is split into 3 levels according to the projection of S . The notation is $J, N \rightarrow J-1, N-1$, where N is the rotational momenta. Beginning from $J = 12$, the transitions are strong and can be observed up to $J = 24$. Additionally, the isotopologue ³³SO is present for $J = 12$ and ³⁴SO for $J = 12$ up to $J = 16$. Weaker SO transitions are present 4_3-1_2 and 5_4-2_3 in Band 1a and 7_7-6_7 up to $12_{12}-11_{12}$. The emission comes predominantly from SMA2, as revealed by the velocity of -7.4 km s^{−1} and the SMA map. The abundance is derived from the isotopologues, but it was necessary to implement a second component which is optically thin.

SiO, Fig. B.1r. Silicon monoxide is viewed as a good tracer for outflows and the $J = 5-4$ transition was observed by Gibb et al. (2007). NGC 6334I shows a bipolar outflow with a red lobe in the northeast and a blue lobe in the southwest direction. With better resolution and additional transitions in the HIFI survey the abundance and spatial position can be given more securely. The $J = 12$ and $J = 13$ transitions in the HIFI data are quite weak. Therefore the parameters were only manually determined. Notable in the SMA data is the narrow -17 km s^{−1} component of the blue lobe and the wide, more diffuse outflow in the red wing. Note that the massive outflows are better traced by CO maps of other groups. Leurini et al. (2006) analyzed the ratios of several CO transitions and derived $1-3 M_{\odot}$ per lobe (dependent on the inclination angle), a dynamical time scale (lobe length divided by v_{\max}) of $2-4 \times 10^3$ yr and a total mass entrainment rate of $\dot{M}_{\text{tot}} = 2-15 \times 10^{-3} M_{\odot}/\text{yr}$.

NH, Fig. B.1d. The ground state transitions $J = 1-0$ of the radical Imidogen with a spin multiplicity of 3 ($^3\Sigma$) and its hyperfine structure are observed. The coupling of the nuclear spin $I = 1$ of ^{14}N with J leads to the splitting into $F = J + I$ components. To fit it properly, a second broad component had to be included. The $J = 1_2-0_1$ transition is blended by the $J = 11-10$ transition of HCN.

SH^+ , Fig. B.1v. Sulfoniumylidene has been detected only recently in Sgr B2(M) (Menten et al. 2011). The 1_2-0_1 transition and its hyperfine structure are observed tentatively above 3σ in absorption and is heavily blended with methanol. The velocity matches that of SMA1. A comparative study of SH^+ absorption lines toward other high-mass star-forming regions has been made by Godard et al. (2012).

CH and CH^+ , Fig. B.1c and B.1b. The hyperfine transitions of methylidene (CH) were analyzed by van der Wiel et al. (2010). We included the blended lines from methanol and find three components: one broad emission from NGC 6334I and two absorption lines from the foreground. The values are in very good agreement with their model I. The two lowest transitions of methylidyne (CH^+) are observed including the isotope ^{13}C (Lis et al., in prep.). The three components are approximately the same as in CH and show the same abundance.

NO, Fig. B.1q. The radical nitrogen monoxide with its doublet structure ($^2\Pi$) is detected in the transitions $^2\Pi_{1/2}$, $J = 5.5-4.5$ (e and f) up to $^2\Pi_{1/2}$, $J = 8.5-7.5$ (e and f). The hyperfine structure is not resolved.

NS, Fig. B.2b. The electronic ground state of nitrogen sulfide is $^2\Pi$, which splits the rotation levels due to the unpaired electron. All transitions from $J = 11$ to 19 are detected. Noticeable is that the fit for the $\Pi_{1/2}$ states is good, but the $\Pi_{3/2}$ states are underexcited. The optical depth is small, so the source size was fixed at $\theta = 3''$.

CN, Fig. B.1l. Due to the weak emission and small number of transitions, there is a large uncertainty of the abundance depending on the velocity width. A width of 4 km s^{-1} would result in an abundance higher by a factor of 2. Here the best fit is taken. The $J = 2_{0,3}-1_{0,2}$ transition at 226.8 GHz in the SMA data is seen in absorption and resembles well the dust emission. The assignment to SMA1 due to $v_{\text{LSR}} = -6.4 \text{ km s}^{-1}$ is confirmed by the map.

CO and CS, Fig. B.1m and B.2a. Carbon monoxide is the most abundant molecule after H_2 in the ISM. In the HIFI spectrum, CO has the most intense lines, which are optically thick. The line profile shows strong outflows with widths of over 80 km s^{-1} and self-absorption. The same is true for carbon monosulfide, so no direct abundance determination is possible. Therefore only the four detected isotopologues of CS and CO are considered, with a focus on C^{18}O and C^{17}O by Ceccarelli et al. (in prep.). For CS, a multi-component approach was tried with two components from the hot cores and one extended emission. But because the first and third component have the same offset velocity, there might be a redundancy.

5.3. Polyatomic linear molecules

N_2H^+ , Fig. B.1g. Four emission lines from $J = 6$ to 10 of diazenylium are detected which exhibit an asymmetric line profile (Ceccarelli et al., in prep.). The intensities could not be fitted properly with a single component. The two component model shows a much better fit, but the values are very uncertain and a narrow line width had to be taken. Russeil et al. (2010) find two components also by observing the hyperfine structure of

the $1_{01}-0_{12}$ line, at -7.3 and -10 km s^{-1} , but with large differences between them concerning T (127 K to 5 K) and Δv (4 to 19 km s^{-1}). Another possibility could be self-absorption.

HNC and HCN, Fig. B.1o and B.1n. Hydrogen isocyanide is an isomer of hydrogen cyanide (HCN). Transitions from $J = 6$ to 10 are present for HNC and $J = 6$ to 12 for HCN. As for N_2H^+ , a two component model was taken. The line profile is changing and affected by self absorption. HCN has a high critical density ($n_{\text{cr}} > 10^9$ for $J = 9-8$ at 797 GHz) and shows an outflow in the red wing of the asymmetric line profile with self-absorption, which indicates an infall. Also vibrationally excited lines ($v_2 = 1, 2$) of HCN are observed. Both aspects, infall and hot HCN, are characteristic for evolved cluster-forming regions (Rolfs et al. 2011). A uniform fit with the isotopologues of HCN is included for the completeness of the survey. For HNC, no other isotopologues are found apart from the HN^{13}C $J = 6-5$ line.

HCO^+ , Fig. B.2p. The rotational transitions of the formyl ion are relatively intense and present up to band 5a (Ceccarelli et al., in prep.). At first, one component including H^{13}CO^+ was tried, but the fitting analysis revealed that either the lower- J or the higher- J transitions are consistent but not all together. This is an indication that either LTE is not satisfied, or several components with different excitation conditions overlap. Assuming two components with different optical depths, the fit could be improved. At higher transitions the component with the higher optical depth dominates.

OCS, Fig. B.2c. Carbonyl sulfide is detected in the rotational transitions from $J = 40$ to 53. Because at first the source size was unknown and no isotopologues are detected in the HIFI spectrum, the abundance was underestimated due to a low opacity ($\tau < 0.3$). With the detection of three isotopologues in the SMA bands the column density was corrected upwards. The emission of O^{13}CS comes predominantly from SMA1.

HC_3N , Fig. B.2t. Cyanoacetylene lines are weak in the HIFI data, but more intense in the SMA data. The fit for SMA1 gives a large line width of 9 km s^{-1} , which is taken for the HIFI model.

5.4. Symmetric top molecules

CH_3CN , Fig. B.2m. Methyl cyanide is considered to be a good tracer of hot cores (Olmi et al. 1993). Rotational transitions from $J = 27$ to 39 with K -components $K = 0-6$ are detected in the HIFI spectrum. Thorwirth et al. (2003) fitted the $J = 12-11$ transition with an unknown source size resulting in $T = 145 \text{ K}$ and $N = 5.2 \times 10^{15} \text{ cm}^{-2}$. The temperature is in good agreement with our value of 154 K. Additionally, weak ($<3\sigma$) bending modes $v_8 = 1$ are observed in band 1a. The small source size and high temperature confirms that methyl cyanide probes the central part of hot cores.

NH_3 , Fig. B.1f. The ground transitions of ammonia are detected mostly in absorption. The overlapping emission and absorption components could not be fitted consistently, and the column density for the emission component associated with NGC 6334I is therefore probably underestimated.

5.5. Inorganic asymmetric molecules

NH_2 , Fig. B.1e. Azanyl is an asymmetric top with a hyperfine structure. Four transitions are detected including the ground state transition. The emission component can be associated with NGC 6334I. As in the case of NH, a broad component ($\Delta v = 30 \text{ km s}^{-1}$) is added to fit the absorption.

H_2Cl^+ and HCl, Fig. B.1u and B.1t. Hydrogen chloride and chloronium were analyzed by Lis et al. (2010). The $J = 2-1$ at 1251 GHz in band 5b of HCl is additionally detected, showing a second component in absorption from the colder and more extended envelope at -10 km s^{-1} with a broad width comparable to H_2Cl^+ . The emission component can be assigned to the center of NGC 6334I. The abundance for this component differs by 50% compared to (Lis et al. 2010) because the source size is not known and the lines are optically thick and affected by the absorption feature. For H_2Cl^+ , the blended lines of dimethylether are included in the fit (blue line). Its velocity and line width leads to the conclusion that it originates from a diffuse foreground molecular cloud.

H_2O^+ , Fig. B.1j. Oxidaniumyl was analyzed by Ossenkopf et al. (2010). Additionally, another weak fine structure transition was observed at 1140 GHz from the $J = 0.5$ level. The derived abundance is above the lower limit given by Ossenkopf et al. (2010). Its radial velocity and line width leads to the conclusion that it is not associated with NGC 6334I.

H_2O , Fig. B.1i. Water was analyzed by Emprechtinger et al. (2010, 2012). It has complex line shapes with many components like optically thick emission and broad outflows, therefore no fit was attempted. Instead, the optically thin isotopes were included to get an estimation of the abundance. Most fits including deuterated water can not reproduce the line strength, so that a non-LTE calculation method is necessary to model them properly.

H_2S , Fig. B.1w. Hydrogen sulfide shows 3 components. The emission one at -6.1 km s^{-1} is affected by two absorption lines, of which one is from the outflow. While the emission peak could be constrained through the isotopologues, the absorptions features could not be reproduced well.

5.6. Organic asymmetric molecules

H_2CO , Fig. B.2l. Formaldehyde is a nearly prolate top and the asymmetry splits every $K \neq 0$ level into doublets. For astrophysical applications, due to its structure this molecule is a good probe of the kinetic temperature and volume density and traces the inner and outer part of hot cores (Mangum & Wootten 1993). Interstellar formaldehyde is formed by hydrogenation of solid CO: $\text{H} + \text{CO} \rightarrow \text{HCO} + \text{H} \rightarrow \text{H}_2\text{CO}$. Another less efficient pathway in the gas phase is $\text{CH}_3 + \text{O} \rightarrow \text{H}_2\text{CO} + \text{H}$. Besides of the weeds, this organic molecule has the most intense lines throughout the HIFI spectrum from band 1a to 4b. The rotational transitions range from $J = 7$ to 15 with $K = 0$ to $K = 6$, and the isotopologues H_2^{13}CO and $\text{H}_2\text{C}^{18}\text{O}$ are observed. The fits with XCLASS show two problems: one component is optically thick, and all the K components could not be described by a single rotation temperature. The $K = 0, 1$ do not match with $K = 2-6$ components consistently. From the SMA data the 3_0-2_0 transition shows that the $K = 0$ value is lower in intensity than the $K = 2$ component. This hints to self absorption in SMA1 or optical thickness.

H_2CCO and H_2CS , Fig. B.2o and B.2f. Ethenone and thioformaldehyde are nearly prolate tops with a spectrum similar to H_2CO . For H_2CCO , only the $K = 1$ components from $J = 24$ to 27 are detected, the other $K = 0$ and $K = 3$ components are included to set an upper limit. Since the transitions are weak and the optical depth is small, θ and Δv needed to be fixed. For H_2CS ranging from $J = 14$ to 23 with $K = 0$ to $K = 5$ transitions, a good fit is obtained where all K components can be described consistently with one excitation temperature.

CH_3OH , Fig. B.2g and B.2h. Methanol is the strongest weed in the HIFI spectrum and accounts for about half of all detected lines. One component is optically thick, which can be derived from the fit of the $^{13}\text{CH}_3\text{OH}$ isotopologue, and the other component accounts for the broad emission. The model fits the observation quite well, except of the most intense lines. As can be seen in the SMA spectra, the XCLASS LTE model cannot reproduce the methanol lines due to a high optical depth. Furthermore, high energy lines $>500 \text{ K}$ are underestimated, so that a 3rd component might be necessary.

NH_2CHO , Fig. B.2s. Formamide is only detected weakly in band 1a. The optical depth is small, and the source size is determined from the SMA map which shows that the emission comes mainly from SMA2.

$\text{C}_2\text{H}_4\text{O}$, Fig. B.2n. Ethylene oxide is surprisingly only detected in SMA2. As for formamide, many weak lines are observed. The source velocity v_{LSR} in the HIFI data is consistent with the general source velocity of SMA2.

HNCO , Fig. B.2r. Isocyanic acid is almost linear, but asymmetric and shows relatively weak emission. The very high temperature and small source size may trace the hot core SMA2.

CH_3OCH_3 , Fig. B.2i. Dimethyl ether has two methyl groups, and similarly to methanol, a splitting into 4 sublevels (AA,EE,AE,EA). It is, after methanol, the organic molecule with the most dense and intense lines in the entire HIFI spectrum. The laboratory measured spectrum from the $^{13}\text{CH}_3\text{OCH}_3$ isotopologue are taken from (Koerber et al. 2010), but are not yet publicly available. The SMA spectra show that the dominant emission arises from SMA1 and that they are weak associations with the outflows.

CH_3OCHO , Fig. B.2j. Methyl formate has the densest spectrum of all detected molecules in the lowest HIFI band 1a (Brouillet et al., in prep.) and the highest number of lines in both SMA bands. The SMA maps show an association with the outflows.

6. Conclusions

In this article, observational data from the HIFI instrument aboard the *Herschel* Space Observatory and from the Submillimeter Array are analyzed for the high-mass star-forming region NGC 6334I. We employ a LTE radiative transfer model to get insights into the physical structure of the source. The molecular spectral line survey gives an overview of the chemical inventory of the molecular species and their excitation, whereas the high-resolution spectrometer allows us tracing the kinematics of the different velocity components. The interferometer maps reveal the source morphology and help to distinguish the embedded cores and resolve their sizes. The results fit well with the overall picture of a high-mass young stellar object and hot molecular cores described by van Dishoeck & Blake (1998).

Two main software tools, XCLASS and MAGIX, are used to demonstrate that modeling and handling of the vast HIFI data sets is possible. These tools allow us to decompose the spectrum into contributions of individual molecules and their isotopologues. By assuming LTE condition and employing the LM fitting algorithm, a fast extraction of information is possible. To take line blends into account, a model for every molecule is produced, which is consistent over multiple HIFI bands. It should be clear that for other excitation conditions than LTE and for optically thick lines like CO this method reaches its limitations. Sophisticated 3-d radiative transfer codes such as

RADMC-3D³ and LIME (Brinch & Hogerheijde 2010) with the possibility of non-LTE calculations will be used in the future to model NGC 6334I into more detail, taking into account the geometry and non-homogeneous distribution of temperature and abundances. With the improved capabilities of ALMA (Atacama Large Millimeter Array), it will be possible to study the molecular chemistry and kinematics on sub-arcsecond scales (<1000 AU) and to follow the fragmentation of protostellar cores with much improved mass sensitivity.

The molecular survey revealed a variety of chemical species, including saturated organic molecules. The astrochemical question is: how can these abundances be explained and how are they related to each other? The next step is therefore a careful comparison of the column densities presented here with predictions of astrochemical models and detailed comparisons with the molecular abundances and abundance ratios in other CHES targets.

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³ <http://www.ita.uni-heidelberg.de/~dullemond/software/radmc-3d/>

Appendix A: SMA spectra and tables

Table A.1. XCLASS fit results from SMA for NGC 6334I.

Source	Species	Θ ($''$)	T_{ex} (K)	N (cm^{-2})	Δv (km s^{-1})	v_{LSR} (km s^{-1})	Notes
SMA1	CN	2	3	0.2×10^{15}	2.2	1.3	
SMA1	CN	2	3	0.8×10^{15}	6.2	-4.6	
SMA2	CN	2	3	0.7×10^{15}	6.9	-0.6	
SMA2	CN	2	3	1.5×10^{15}	3.2	-4.5	
SMA1	DCN	3.5	359	1.6×10^{15}	5.8	-4.9	2
SMA2	DCN	5.0	146	0.6×10^{15}	6	-7.8	2
SMA1	SiO	1.7	40	1.0×10^{15}	8	-3.2	2
SMA2	SiO	2.0	40	1.0×10^{15}	6	-5.2	2
SMA1	H ₂ S	3	62	1.4×10^{17}	5	-4	2, 3
SMA2	H ₂ S	3	147	0.4×10^{17}	6.5	-6	2, 3
SMA1	CS	2.4	40	1.6×10^{18}	5.1	-3.2	2, 3, 5
SMA2	CS	2.2	32	6.6×10^{17}	2.7	-6.4	2, 3, 5
SMA1	OCS	2.5	137	2.3×10^{18}	4	-5.6	1
SMA2	OCS	2.5	126	1.1×10^{18}	4	-7	1
SMA1	CH ₃ OH	4	78	2.4×10^{19}	3.6	-5.2	
SMA2	CH ₃ OH	4	104	2.0×10^{19}	3.0	-7.5	
SMA1	CH ₃ OCH ₃	3.0	78	2.3×10^{18}	4.9	-5.6	
SMA2	CH ₃ OCH ₃	3.0	98	0.7×10^{18}	5.4	-7.7	
SMA1	CH ₃ OCHO	2.3	104	7.7×10^{17}	4.6	-6.1	
SMA2	CH ₃ OCHO	3	115	9.6×10^{17}	3.8	-7.6	
SMA1	H ₂ CO	3.0	104	1.6×10^{17}	5.7	-5.5	
SMA2	H ₂ CO	3.0	138	1.3×10^{17}	4.6	-7.5	
SMA1	NH ₂ CHO	2	25	4.7×10^{15}	5.2	-4.7	2, 3
SMA2	NH ₂ CHO	2	45	5.0×10^{15}	4.0	-6.9	2, 3
SMA1	HC ₃ N	2.8	208	4.2×10^{15}	9.4	-5.7	
SMA2	HC ₃ N	2.4	223	1.4×10^{15}	5.8	-6.6	
SMA1	C ₂ H ₃ CN	2.0	147	5.2×10^{15}	6.0	-6.3	
SMA1	C ₂ H ₅ CN	2.2	280	4.2×10^{16}	6.6	-5.5	
SMA2	C ₂ H ₅ CN	1.0	123	1.6×10^{16}	4.1	-6.7	
SMA1	CH ₃ CH ₃ CO	2.6	73	2.1×10^{16}	5.6	-5.4	
SMA2	CH ₃ CH ₃ CO	2.6	159	2.9×10^{16}	4.8	-7.4	
SMA2	HCOOH	2.0	258	5.5×10^{16}	4.3	-9	2
SMA2	(CH ₂ OH) ₂	2.2	217	2.6×10^{16}	3.4	-8.2	
SMA2	C ₂ H ₅ OH	2.5	77	1.4×10^{17}	4.9	-7.9	
SMA2	C ₂ H ₄ O	2.1	335	6.7×10^{16}	5	-6.8	
SMA2	SO	2.5	70	1.7×10^{17}	4.2	-8.7	2, 6

Notes. (1) Including only isotopologues ^{18}OCS , O^{13}CS , OC^{33}S . (2) 1 or few lines. (3) Blended. (4) Including $^{13}\text{CH}_3\text{OH}$. (5) Including only $^{13}\text{C}^{34}\text{S}$. (6) Including only ^{33}SO .

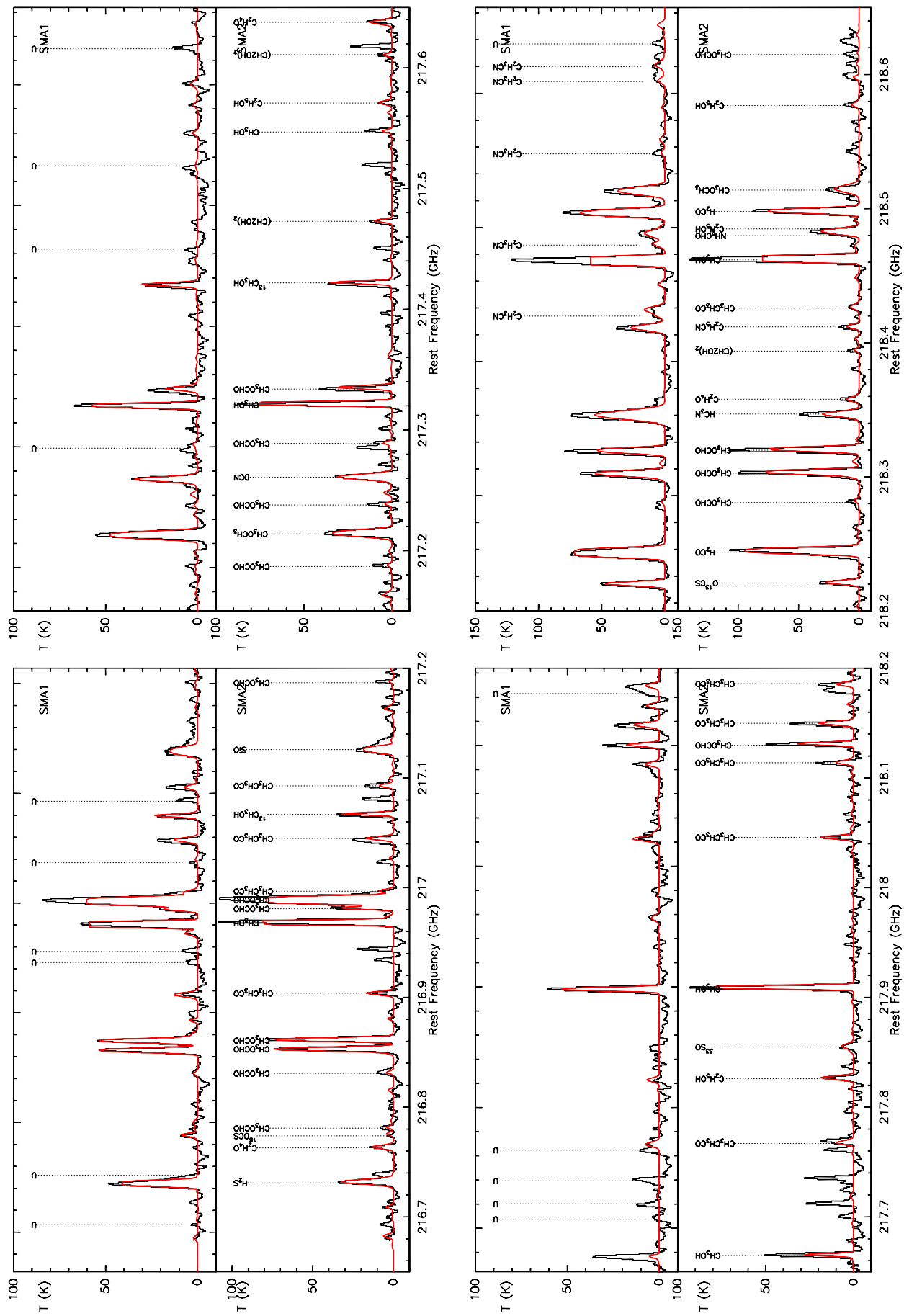


Fig. A.1. LSB Spectra (in steps of 500 MHz) of the two hot cores SMA1 (*upper*) and SMA2 (*lower*) of NGC 6334I, observed with the Submillimeter Array. Overlaid in red are the XCLASS fits. For more detailed information about each identified spectral line, see Table A.2.

Fig. A.1. continued.

Table A.2. List of identified lines in the SMA bands of NGC 6334I in order of increasing rest frequency.

Molecule	Transition	LSB Frequency (MHz)	Molecule	Transition	USB Frequency (MHz)
H ₂ S	2 2 0–2 1 1	216 710.437	C ₂ H ₅ OH	10 2 9 2–9 1 8 2	226 661.701
C ₂ H ₄ O	16 5 11–16 4 12	216 741.468	CN	2 0 2 2–1 0 1 1	226 679.311
C ₂ H ₄ O	16 6 11–16 5 12	216 741.712	CH ₃ CH ₃ CO, $v = 1$	20 3 17 0–19 4 16 1	226 692.48
¹⁸ OCS	19 –18	216 753.488	CH ₃ CH ₃ CO, $v = 1$	20 4 17 0–19 3 16 1	226 692.48
CH ₃ OCHO	32 6 27 1–32 5 28 1	216 758.969	CH ₃ OCHO	20 2 19 1–19 2 18 1	226 713.06
CH ₃ OCHO	32 6 27 0–32 5 28 0	216 809.014	CH ₃ OCHO	20 2 19 0–19 2 18 0	226 718.688
CH ₃ OCHO	18 2 16 2–17 2 15 2	216 830.197	¹³ CH ₃ OH	12 5 8 0–13 4 9 0	226 716.392
CH ₃ OCHO	18 2 16 0–17 2 15 0	216 838.889	CH ₃ CH ₃ CO, $v = 0$	20 4 17 1–19 3 16 1	226 784.639
CH ₃ CH ₃ CO, $v = 1$	19 4 16 0–18 4 15 1	216 881.337	CH ₃ CH ₃ CO, $v = 0$	20 3 17 1–19 4 16 1	226 784.639
CH ₃ CH ₃ CO, $v = 1$	19 3 16 0–18 4 15 1	216 881.337	CH ₃ CH ₃ CO, $v = 0$	20 4 17 1–19 3 16 2	226 784.67
CH ₃ CH ₃ CO, $v = 1$	19 4 16 0–18 3 15 1	216 881.337	CH ₃ CH ₃ CO, $v = 0$	20 3 17 1–19 4 16 2	226 784.67
CH ₃ CH ₃ CO, $v = 1$	19 3 16 0–18 3 15 1	216 881.337	CH ₃ OCHO	20 1 19 2–19 1 18 2	226 773.13
CH ₃ OH	5 1 0 0–4 2 0 0	216 945.521	CH ₃ OCHO	20 1 19 0–19 1 18 0	226 778.786
CH ₃ OCHO	17 3 14 5–16 3 13 5	216 958.834	CH ₃ CH ₃ CO, $v = 0$	20 4 17 0–19 3 16 1	226 832.055
CH ₃ OCHO	20 0 20 2–19 1 19 1	216 962.989	CH ₃ CH ₃ CO, $v = 0$	20 3 17 0–19 4 16 1	226 832.055
CH ₃ CH ₃ CO, $v = 0$	19 4 16 1–18 3 15 1	216 974.426	CH ₃ OCHO	20 2 19 1–19 1 18 2	226 856.872
CH ₃ CH ₃ CO, $v = 0$	19 3 16 1–18 4 15 1	216 974.426	CH ₃ OCHO	20 2 19 0–19 1 18 0	226 862.257
CH ₃ CH ₃ CO, $v = 0$	19 4 16 1–18 3 15 2	216 974.464	CN	2 0 3 3–1 0 2 2	226 874.191
CH ₃ CH ₃ CO, $v = 0$	19 3 16 1–18 4 15 2	216 974.464	CN	2 0 3 4–1 0 2 3	226 874.781
CH ₃ CH ₃ CO, $v = 0$	19 4 16 0–18 3 15 1	217 022.509	CH ₃ CH ₃ CO, $v = 0$	20 4 17 0–19 3 16 0	226 879.393
CH ₃ CH ₃ CO, $v = 0$	19 3 16 0–18 4 15 1	217 022.509	CH ₃ CH ₃ CO, $v = 0$	20 3 17 0–19 4 16 0	226 879.393
¹³ CH ₃ OH	14 1 13 0–13 2 12 0	217 044.616	CN	2 0 3 2–1 0 2 2	226 887.42
CH ₃ CH ₃ CO, $v = 0$	19 4 16 0–18 3 15 0	217 070.504	CN	2 0 3 3–1 0 2 3	226 892.128
CH ₃ CH ₃ CO, $v = 0$	19 3 16 0–18 4 15 0	217 070.504	CH ₃ OCHO	6 6 1 4–5 5 1 4	226 913.736
CH ₃ OCHO	30 4 26 0–30 4 27 0	217 077.079	CH ₃ OH	14–2 0 0–13 3 0 0	226 939.41
SiO	5–4	217 104.98	CH ₃ OCHO	25 9 17 1–25 8 18 1	226 974.039
CH ₃ OCHO	35 9 27 1–35 8 28 1	217 165.366	CH ₃ OCHO	19 2 17 2–18 2 16 2	227 019.55
CH ₃ OCH ₃	22 4 19 3–22 3 20 3	217 189.668	CH ₃ OCHO	19 2 17 0–18 2 16 0	227 028.121
CH ₃ OCHO	32 9 24 0–32 8 25 0	217 215.847	CH ₃ CH ₃ CO, $v = 1$	20 3 17 1–19 4 16 1	227 032.073
DCN	3 2 –2 2	217 238.631	CH ₃ CH ₃ CO, $v = 1$	20 4 17 1–19 3 16 1	227 032.073
CH ₃ OCHO	30 4 26 0–30 3 27 0	217 266.549	CH ₃ CH ₃ CO, $v = 1$	20 4 17 1–19 3 16 2	227 043.496
CH ₃ OH	6 1–1 1–7 2–1 1	217 299.205	CH ₃ CH ₃ CO, $v = 1$	20 3 17 1–19 4 16 2	227 043.496
CH ₃ OCHO	17 4 13 3–16 4 12 3	217 312.626	CH ₃ OH	21 1 0 0–21 0 0 0	227 094.747
¹³ CH ₃ OH	10 2 8 0–9 3 7 0	217 399.55	CH ₃ OCHO	25 9 16 2–25 8 18 1	227 106.67
(CH ₂ OH) ₂	24 1 24 0–23 1 23 1	217 449.995	CH ₃ OH	12–1 0 0–11 2 0 0	227 229.511
(CH ₂ OH) ₂	24 0 24 0–23 0 23 1	217 450.27	C ₂ H ₅ OH	13 3 10 1–12 3 9 1	227 294.752
CH ₃ OH	16 1 0 0–15 3 0 0	217 525.002	¹³ C ³⁴ S	5–4	227 300.506
C ₂ H ₅ OH	5 1 4 1–4 0 4 0	217 548.152	HC ₃ N, $v = 0$	25 –24	227 418.905
C ₂ H ₅ OH	25 3 22 0–24 4 20 1	217 549.328	CH ₃ CH ₃ CO, $v = 1$	21 3 19 0–20 3 18 0	227 560.862
(CH ₂ OH) ₂	21 2 19 1–20 2 18 0	217 587.548	CH ₃ CH ₃ CO, $v = 1$	21 2 19 0–20 2 18 0	227 560.862
C ₂ H ₄ O	15 4 11–15 3 12	217 615.673	CH ₃ OCHO	21 0 21 2–20 0 20 2	227 561.741
C ₂ H ₄ O	15 5 11–15 4 12	217 615.72	(CH ₂ OH) ₂	22 4 18 0–21 4 17 1	227 587.074
CH ₃ OH	15 6 1 1–16 5 1 1	217 642.678	CH ₃ OCHO	18 3 15 3–17 3 14 3	227 599.261
CH ₃ CH ₃ CO, $v = 1$	20 2 18 0–19 2 17 0	217 744.747	NH ₂ CHO	11 0 11–10 0 10	227 606.176
CH ₃ CH ₃ CO, $v = 1$	20 3 18 0–19 3 17 0	217 744.747	C ₂ H ₅ OH	18 5 13 2–18 4 14 2	227 606.079
C ₂ H ₅ OH	5 3 3 2–4 2 2 2	217 803.689	C ₂ H ₅ CN, $v = 1$	26 1 26 1–25 1 25 1	227 624.735
³³ SO	5 6 8–4 5 7	217 832.642	CH ₃ OCH ₃	47 11 36 1–46 12 35 1	227 668.331
CH ₃ OH	20 1 0 0–20 0 0 0	217 886.504	CH ₃ OCH ₃	47 11 37 3–46 12 34 3	227 675.733
CH ₃ CH ₃ CO, $v = 1$	20 2 18 0–19 3 17 1	218 023.182	C ₂ H ₅ CN	25 3 22–24 3 21	227 780.972
CH ₃ CH ₃ CO, $v = 1$	20 3 18 0–19 2 17 1	218 023.182	CH ₃ OH	16 1 1 0–15 2 1 0	227 814.528
CH ₃ CH ₃ CO, $v = 0$	20 3 18 1–19 2 17 2	218 091.448	CH ₃ CH ₃ CO, $v = 1$	21 3 19 0–20 3 18 1	227 836.273
CH ₃ CH ₃ CO, $v = 0$	20 2 18 1–19 3 17 2	218 091.448	CH ₃ CH ₃ CO, $v = 1$	21 2 19 0–20 2 18 1	227 836.273
CH ₃ OCHO	17 4 13 5–16 4 12 5	218 108.438	C ₂ H ₃ CN	24 7 18–23 7 17	227 897.607
CH ₃ CH ₃ CO, $v = 0$	20 2 18 0–19 3 17 1	218 127.207	C ₂ H ₃ CN	24 7 17–23 7 16	227 897.608
CH ₃ CH ₃ CO, $v = 0$	20 3 18 0–19 2 17 1	218 127.207	CH ₃ CH ₃ CO, $v = 0$	21 3 19 1–20 3 18 1	227 903.922
CH ₃ CH ₃ CO, $v = 0$	20 3 18 0–19 2 17 0	218 162.929	CH ₃ CH ₃ CO, $v = 0$	21 2 19 1–20 2 18 1	227 903.922
CH ₃ CH ₃ CO, $v = 0$	20 2 18 0–19 3 17 0	218 162.929	CH ₃ CH ₃ CO, $v = 0$	21 3 19 1–20 3 18 2	227 903.953
O ¹³ CS	18 –17	218 198.998	CH ₃ CH ₃ CO, $v = 0$	21 2 19 1–20 2 18 2	227 903.953
H ₂ CO	3 0 3–2 0 2	218 222.192	C ₂ H ₃ CN	24 6 19–23 6 18	227 906.683
CH ₃ OCHO	31 9 23 0–31 8 24 0	218 259.594	C ₂ H ₃ CN	24 6 18–23 6 17	227 906.709
CH ₃ OCHO	17 3 14 2–16 3 13 2	218 280.9	CH ₃ CH ₃ CO, $v = 0$	21 2 19 0–20 2 18 1	227 939.374
CH ₃ OCHO	17 3 14 0–16 3 13 0	218 297.89	CH ₃ CH ₃ CO, $v = 0$	21 2 19 0–20 3 18 1	227 939.374
HC ₃ N, $v = 0$	24 –23	218 324.723	CH ₃ CH ₃ CO, $v = 0$	21 3 19 0–20 2 18 1	227 939.374
C ₂ H ₄ O	14 3 11–14 2 12	218 335.626	CH ₃ CH ₃ CO, $v = 0$	21 3 19 0–20 3 18 1	227 939.374

Table A.2. continued.

LSB			USB		
Molecule	Transition	Frequency (MHz)	Molecule	Transition	Frequency (MHz)
C ₂ H ₄ O	14 4 11–14 3 12	218 335.634	CH ₃ OCHO	24 9 15 0–24 8 16 0	227 954.452
(CH ₂ OH) ₂	22 4 19 0–21 4 18 1	218 371.495	C ₂ H ₃ CN	24 5 20–23 5 19	227 966.032
C ₂ H ₅ CN	24 3 21–23 3 20	218 389.97	C ₂ H ₃ CN	24 5 19–23 5 18	227 967.589
C ₂ H ₃ CN	23 7 17–22 7 16	218 398.555	CH ₃ CH ₃ CO, $v = 0$	21 3 19 0–20 3 18 0	227 974.758
C ₂ H ₃ CN	23 7 17–22 7 16	218 398.555	CH ₃ CH ₃ CO, $v = 0$	21 2 19 0–20 2 18 0	227 974.758
C ₂ H ₃ CN	23 7 16–22 7 15	218 398.555	HC ₃ N, $v_7=1$	25 1 –24–1	227 977.277
C ₂ H ₃ CN	23 6 18–22 6 17	218 402.435	CH ₃ OCHO	31 4 27 2–31 3 28 2	227 986.255
C ₂ H ₃ CN	23 6 17–22 6 16	218 402.451	CH ₃ OCHO	24 9 15 2–24 8 16 2	227 994.545
CH ₃ CH ₃ CO, $v = 1$	11 11 1 0–10 10 0 0	218 402.853	C ₂ H ₅ OH	13 1 12 1–12 1 11 1	227 891.911
CH ₃ CH ₃ CO, $v = 1$	11 11 0 0–10 10 1 0	218 404.816	C ₂ H ₅ OH	13 3 10 0–12 3 9 0	228 029.05
CH ₃ OH	4 2 0 0–3 1 0 0	218 440.063	CH ₃ OCHO	31 4 27 0–31 3 28 0	228 057.899
C ₂ H ₃ CN	23 5 19–22 5 18	218 451.297	C ₂ H ₃ CN	24 11 13–23 11 12	228 087.245
C ₂ H ₃ CN	23 5 18–22 5 17	218 452.357	C ₂ H ₃ CN	24 11 14–23 11 13	228 087.245
NH ₂ CHO	10 1 9–9 1 8	218 459.653	C ₂ H ₃ CN	24 4 21–23 4 20	228 104.614
C ₂ H ₅ OH	5 3 2 2–4 2 3 2	218 461.226	OC ³³ S	19 –18	228 151.943
H ₂ CO	3 2 2–2 2 1	218 475.632	C ₂ H ₃ CN	24 4 20–23 4 19	228 160.305
CH ₃ OCH ₃	23 3 21 1–23 2 22 1	218 491.914	CH ₃ OCHO	24 9 16 1–24 8 17 1	228 205.832
CH ₃ OCH ₃	23 3 21 0–23 2 22 0	218 494.39	CH ₃ OCHO	18 3 15 5–17 3 14 5	228 211.291
C ₂ H ₃ CN	23 10 13–22 10 12	218 519.997	CH ₃ OCHO	24 9 16 0–24 8 17 0	228 270.535
C ₂ H ₃ CN	23 10 14–22 10 13	218 519.997	HC ₃ N, $v_7=1$	25–1–24 1	228 303.174
C ₂ H ₅ OH	21 5 16 2–21 4 17 2	218 554.499	CH ₃ OCH ₃	26 3 24 0–25 4 21 0	228 421.737
C ₂ H ₃ CN	23 4 20–22 4 19	218 573.646	CH ₃ OCH ₃	26 3 24 1–25 4 21 1	228 422.878
C ₂ H ₃ CN	23 3 21–22 3 20	218 585.072	CH ₃ OH	21 0 0 0–20 3 0 0	228 467.902
CH ₃ OCHO	27 7 21 0–27 5 22 0	218 593.341	C ₂ H ₅ CN	25 2 23–24 2 22	228 483.136
			HCOOH	10 2 8–9 2 7	228 544.168

Appendix B: HIFI fitted curves of all molecules

Examples of selected, fitted spectra for each molecule. In red is the model, in blue are mostly annotations or blended lines. If only few transitions of a molecule are present, they are all shown in the plot. For weeds or molecules with many lines, no quantum numbers are given and only selected frequency ranges over different bands are presented. Heavy blended lines are excluded for simplicity. All intensities are given in units of Kelvin. The quantum numbers refer to upper state – lower state. For absorption lines, this sequence should be read reversed.

The entire observed and simulated spectrum for each HIFI band will be made available online in the CLASS format on the CHESS website (<http://www-laog.obs.ujf-grenoble.fr/heberges/hs3f/>), as soon as the HIPE 8 reprocessed data are available. Additionally, the modeled spectrum for each molecule (including isotopologues and vibrationally excited states) can be loaded. It should be noticed that for the simulated spectrum of only one molecule, the blending by other molecules is not taken into account. This results in a deviation for the fits of absorption lines when a nearby emission line contributes to the continuum. In XCLASS, at first all the emission components are calculated and then the absorption components.

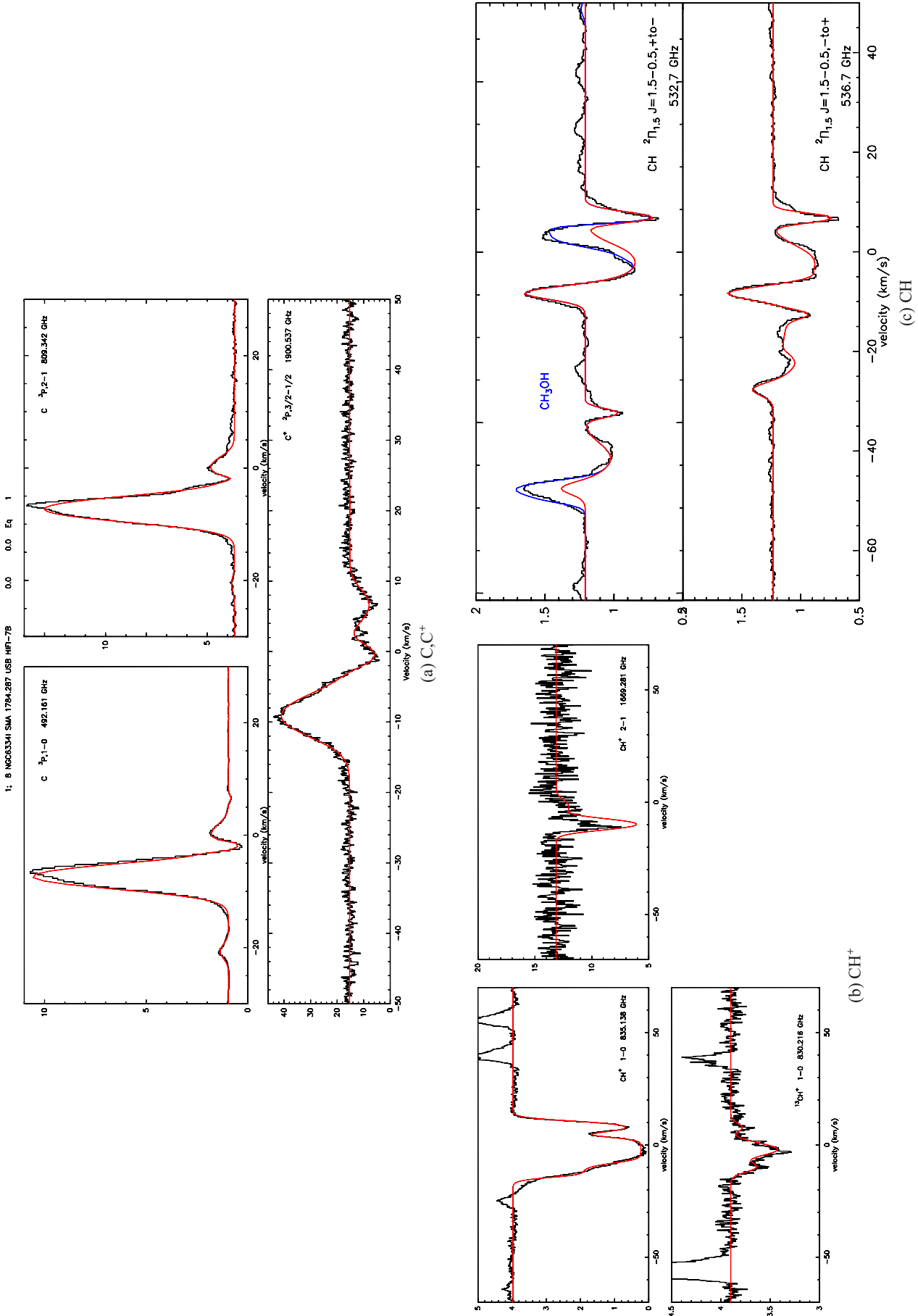


Fig B.1. XCLASS fits of chemical species in the HIFI spectrum of NGC 6334I.

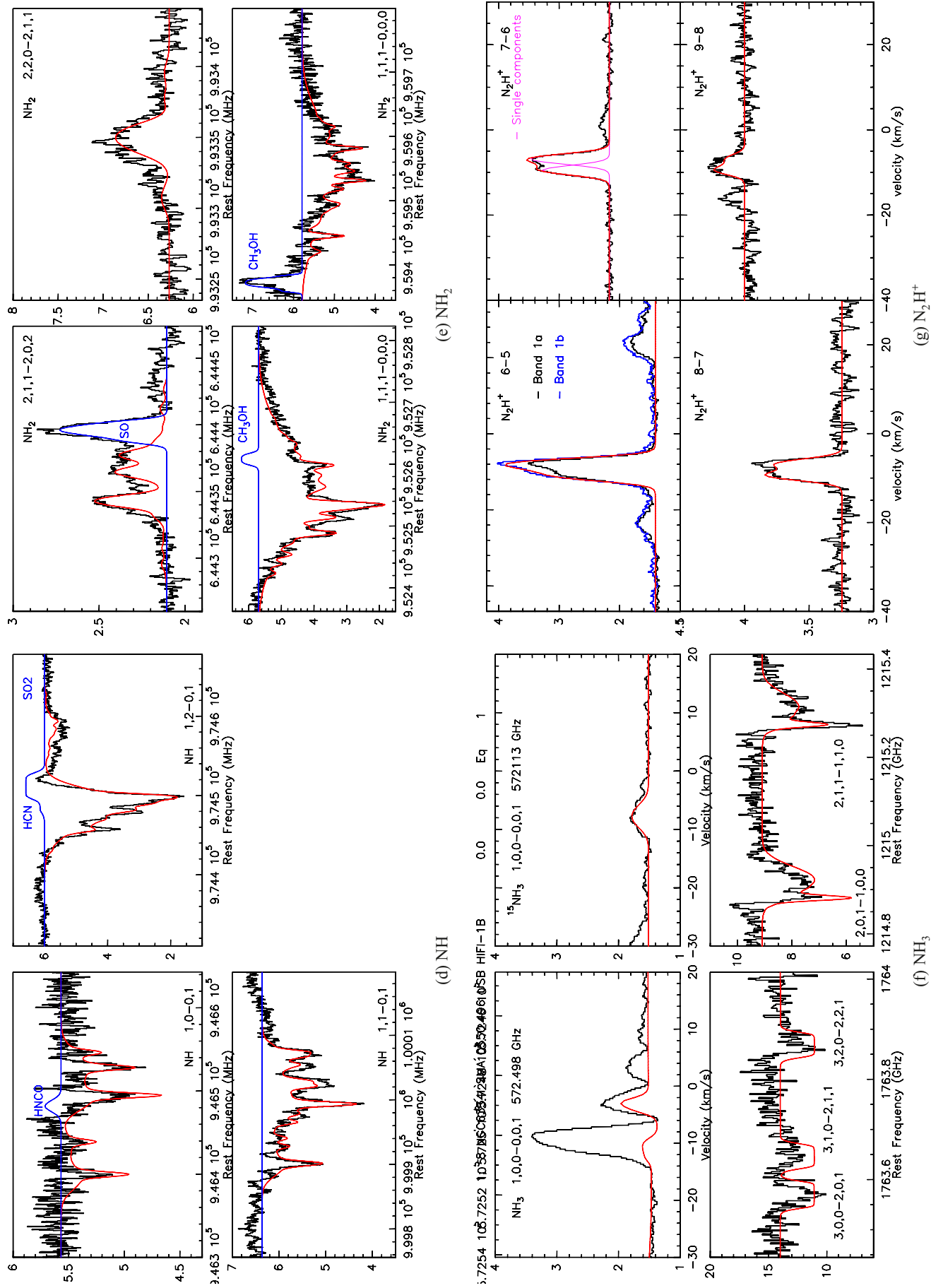


Fig. B.1. continued.

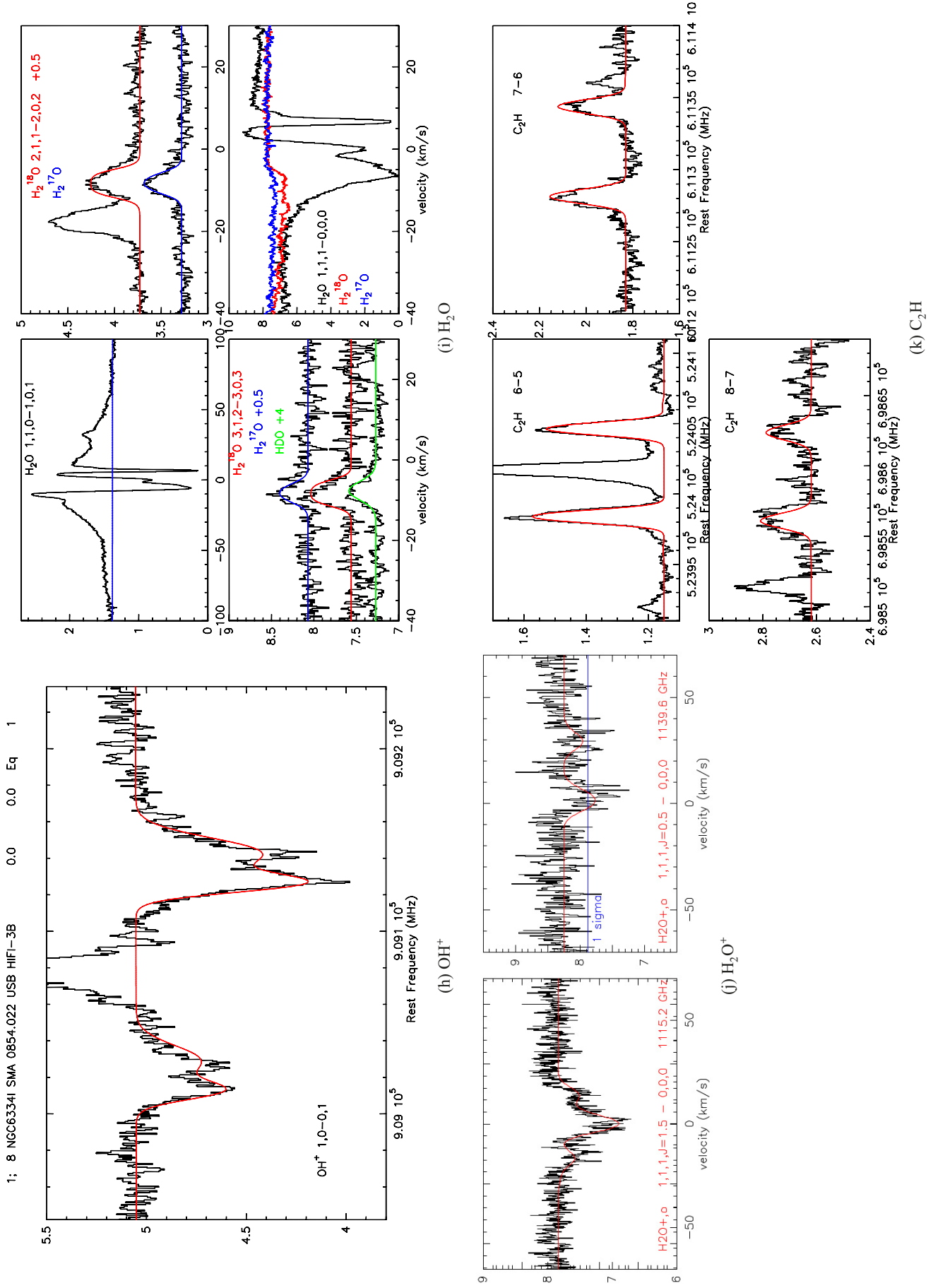


Fig. B.1.1. continued.

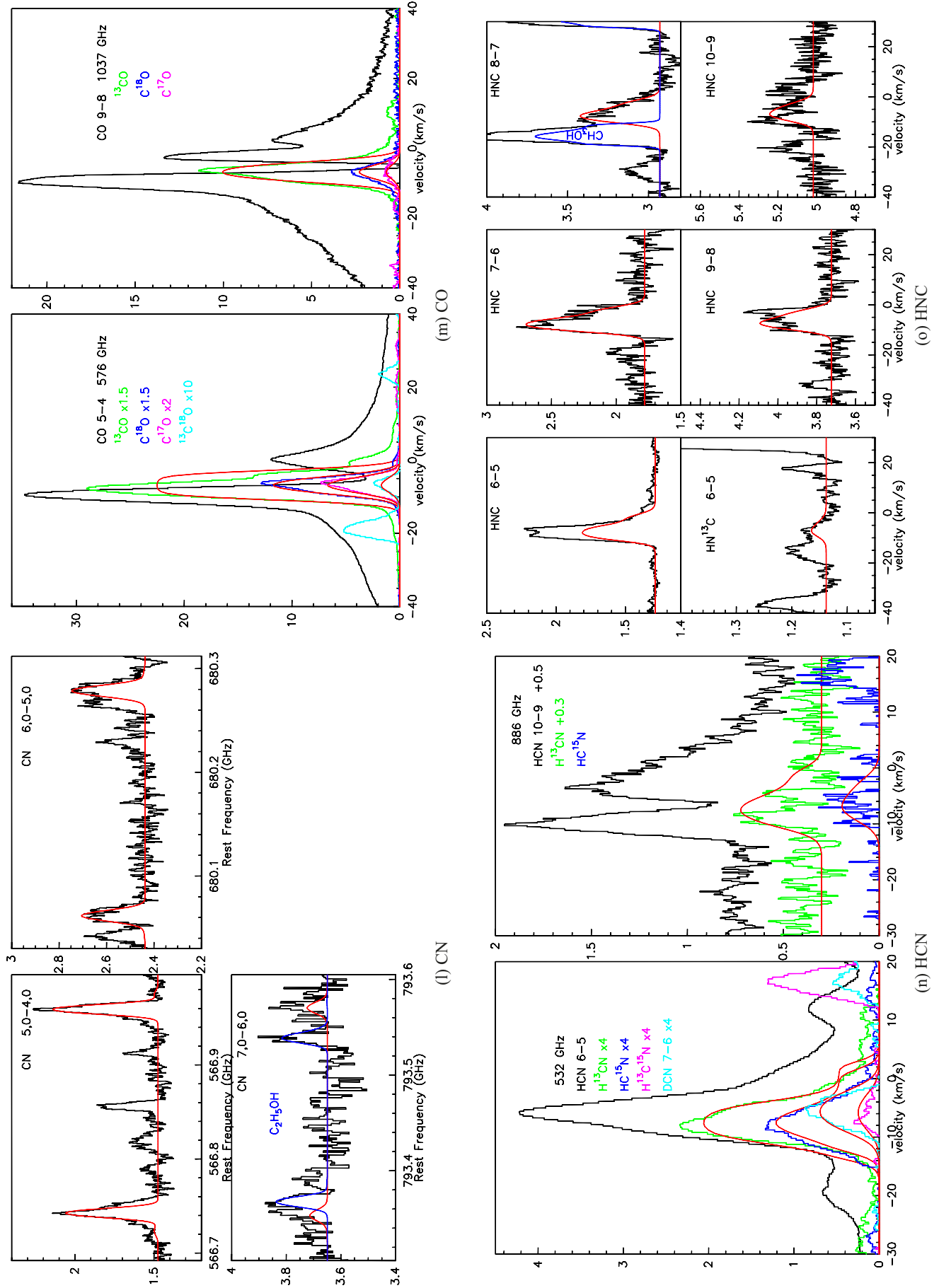


Fig. B.1. continued.

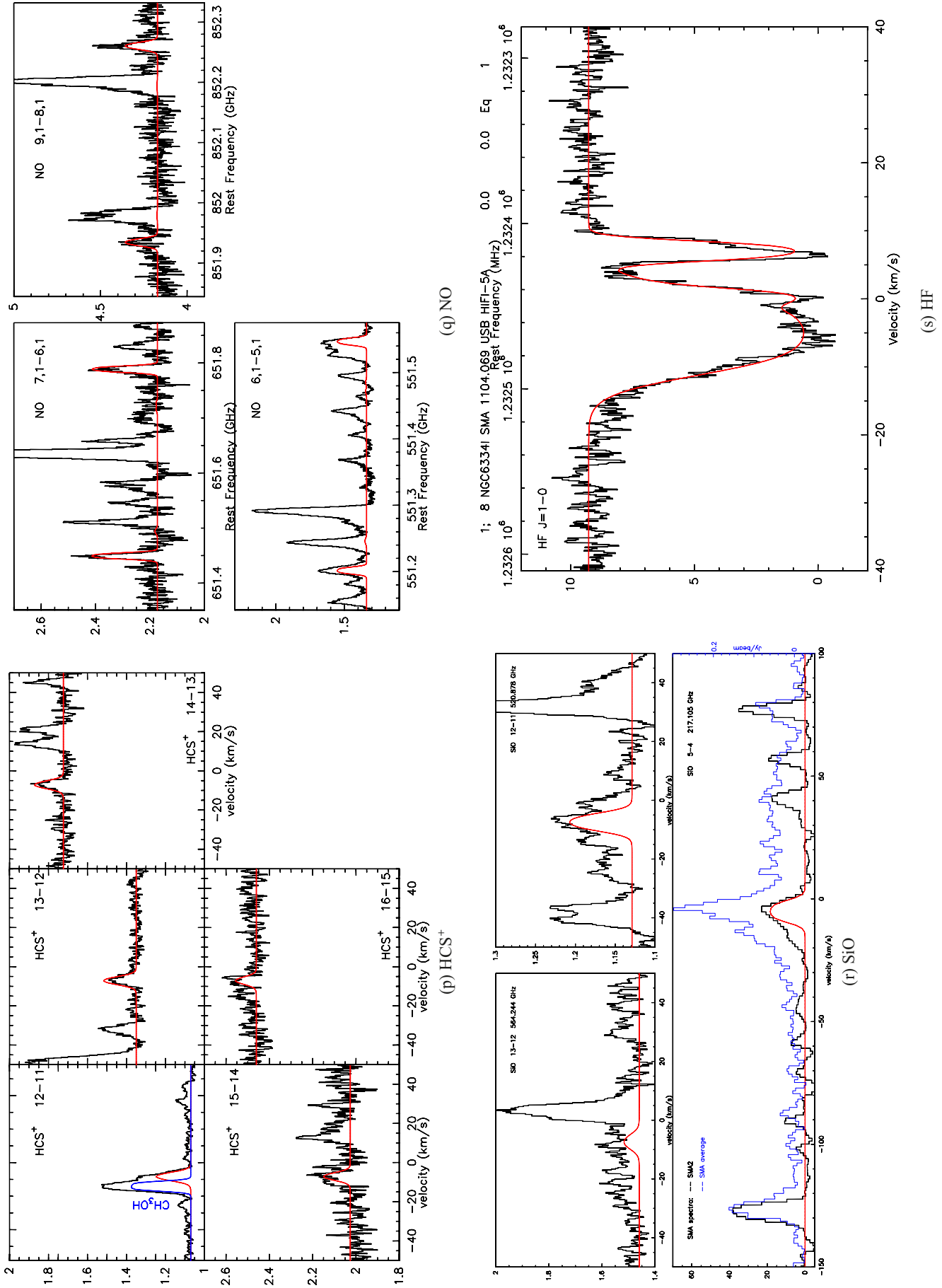


Fig. B.1. continued.

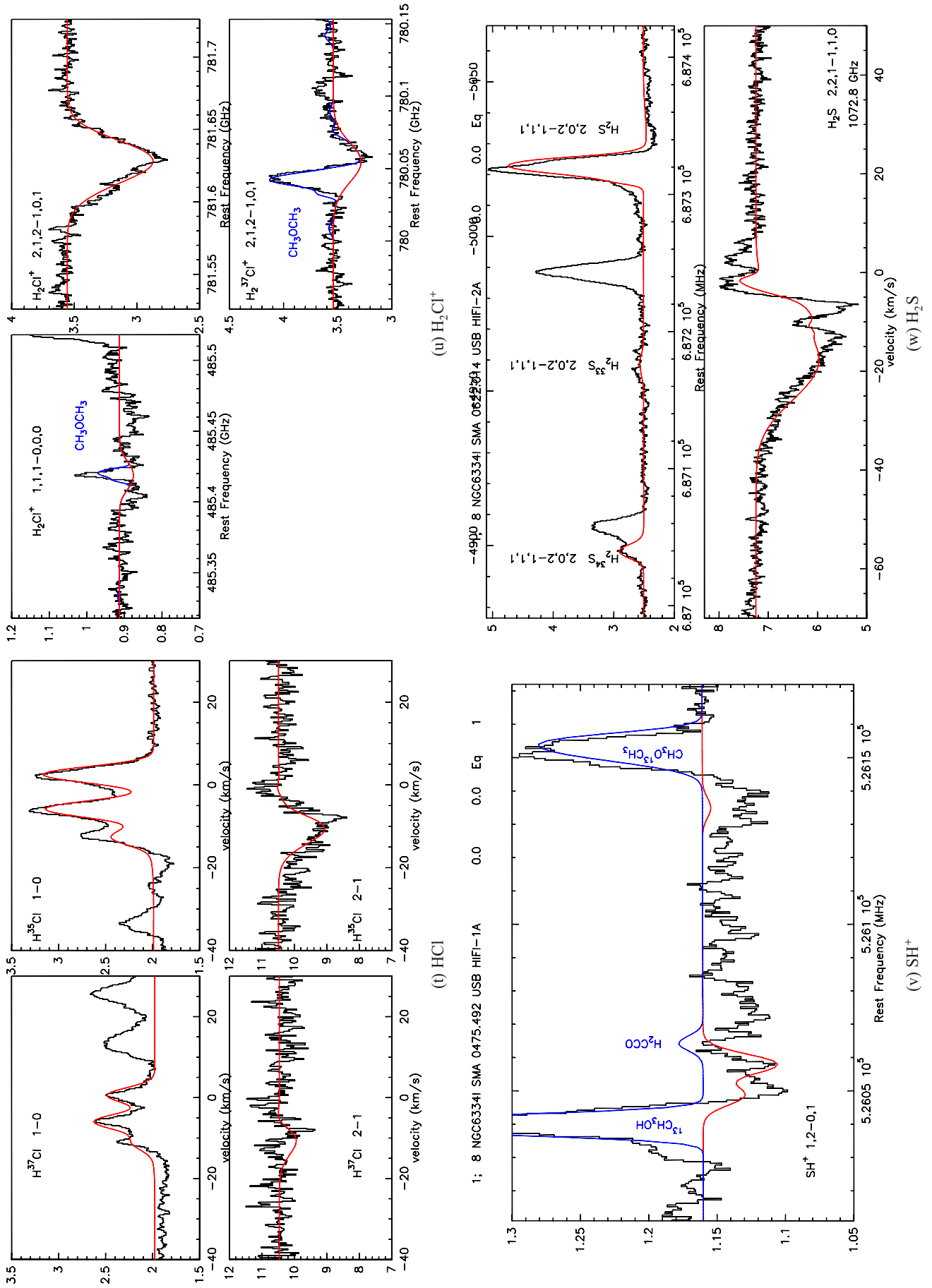


Fig. B.1. continued.

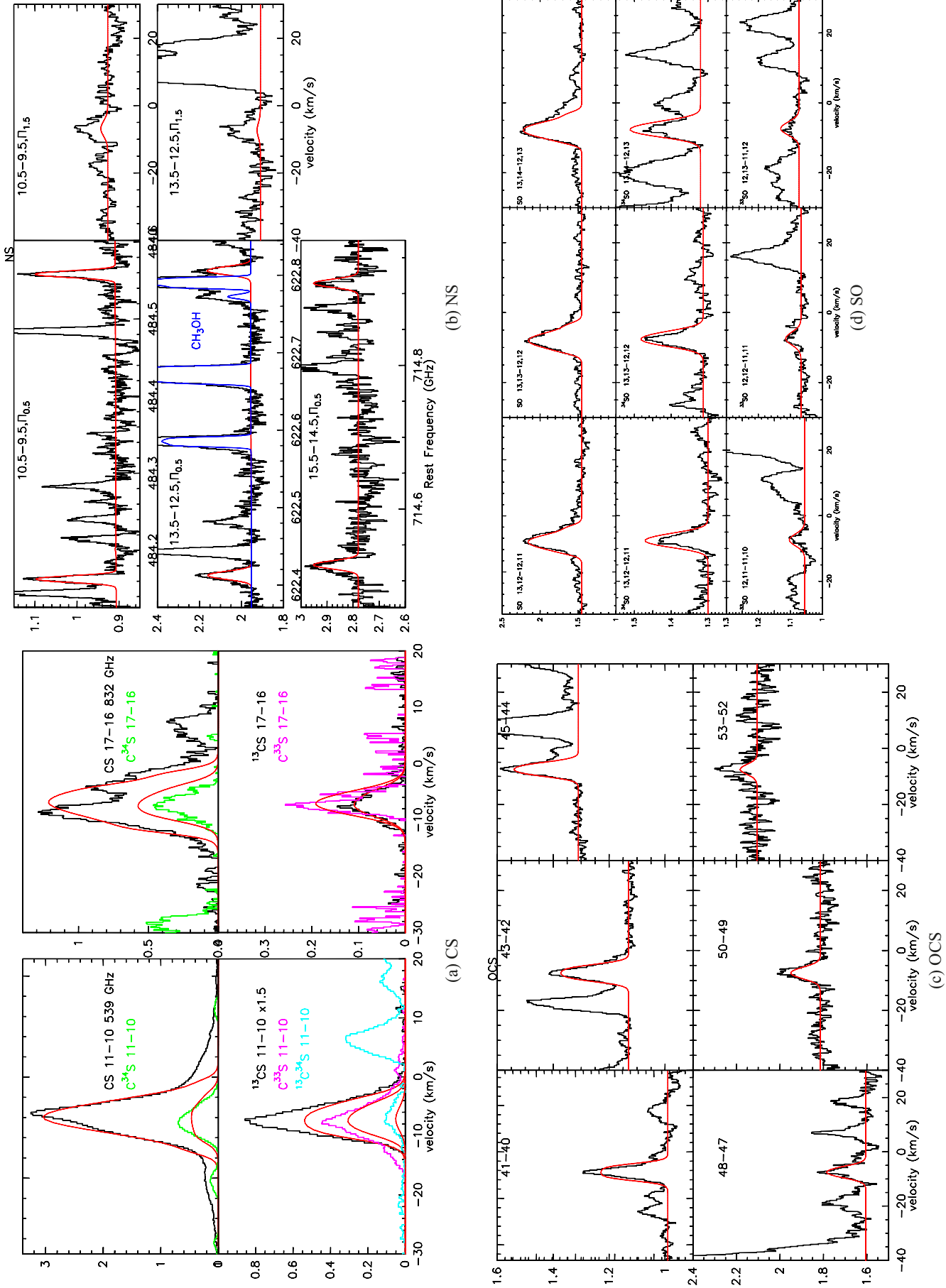


Fig. B.2. XCLASS fits of chemical species in the HIFI spectrum of NGC 6334I, continued.

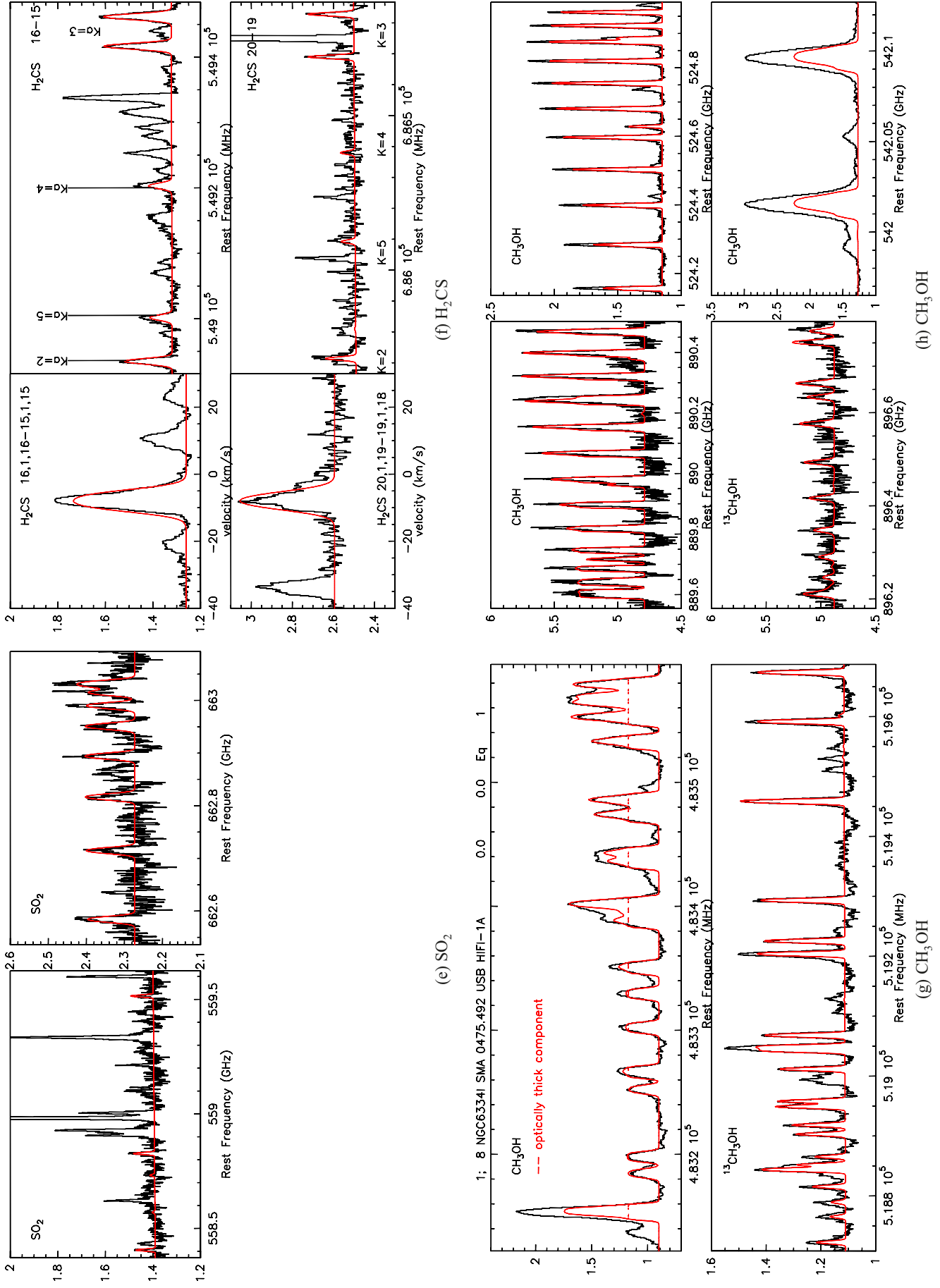


Fig. B.2. continued.

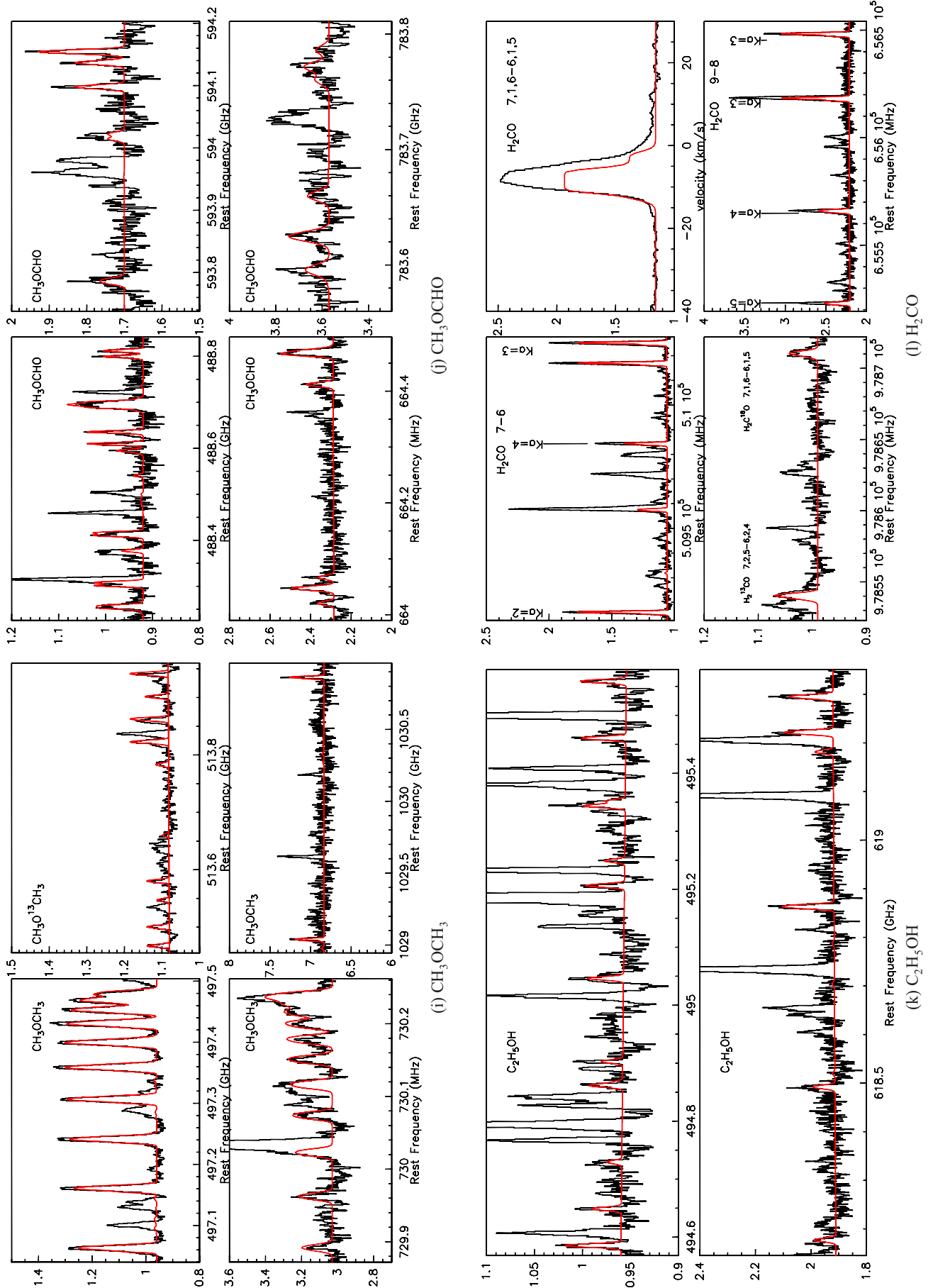


Fig. B.2. continued.

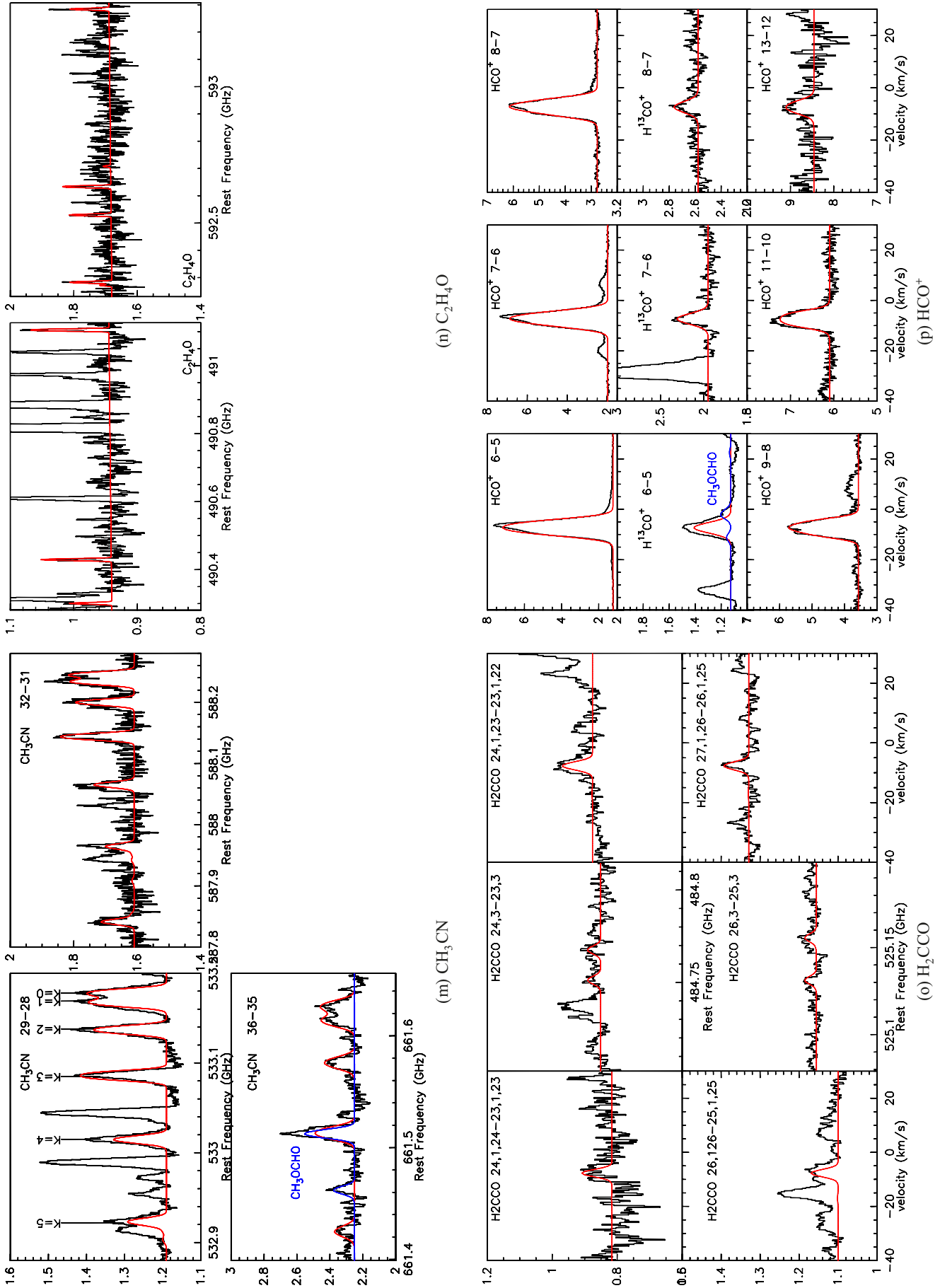


Fig. B.2. continued.

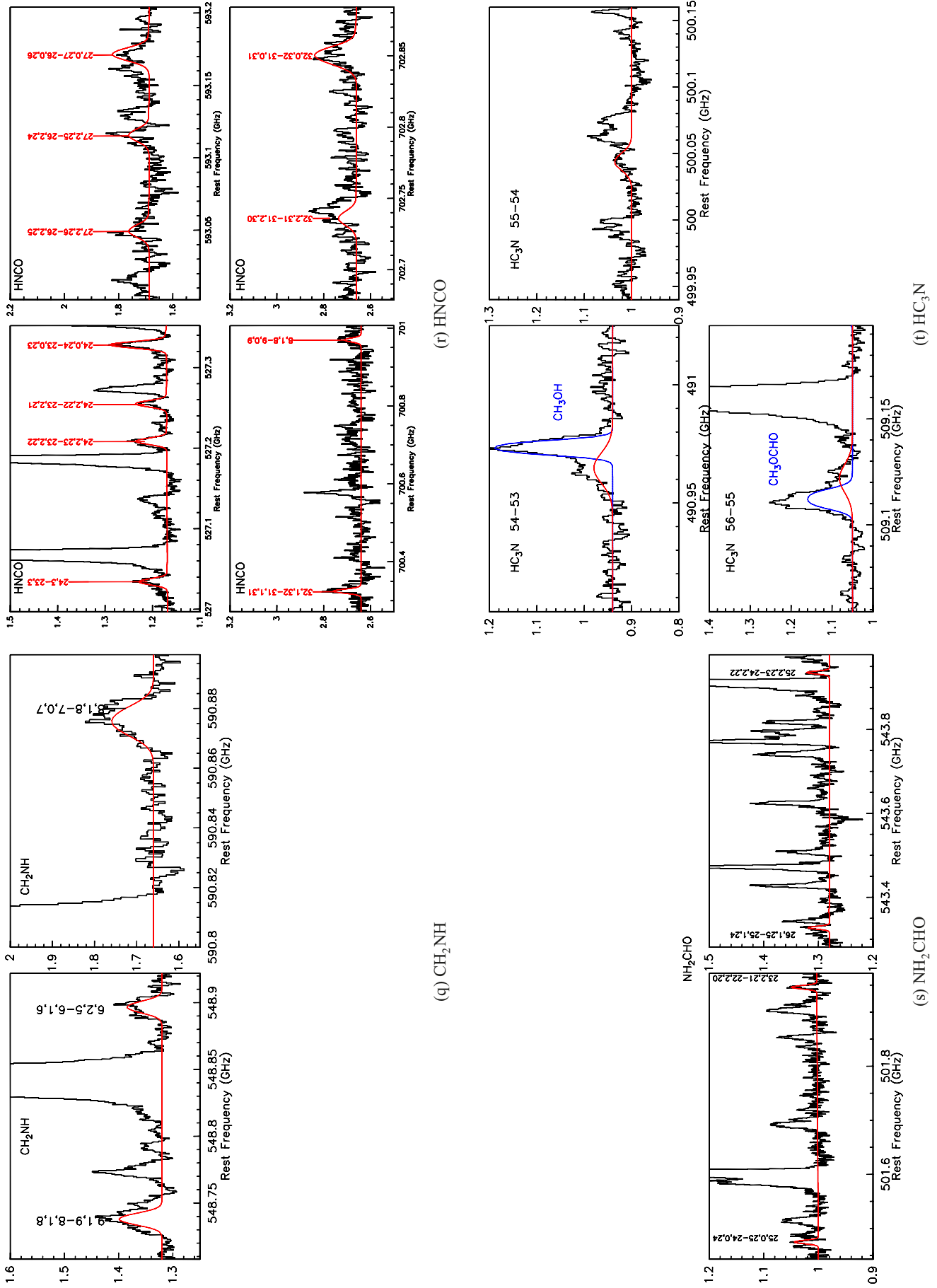


Fig. B.2. continued.